An aerial photograph of Catania, Sicily, Italy. The city is densely packed with buildings, extending from the coast inland. In the foreground, a large harbor is visible with several ships and a long pier. The background features the prominent, snow-capped Mount Etna under a clear blue sky.

# Breakdown of the Universality of Glasses at Ultralow Temperatures – Interplay of Nuclear Spins and Atomic Tunneling Systems

Catania 1999





# Comment Paper by Clare and Tony 1988

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## Low Temperature Properties of Amorphous Materials: Through a Glass Darkly

A wide variety of amorphous materials exhibit similar behavior in their thermal properties. Examples of such universal features include a linear specific heat and a  $T^2$  thermal conductivity below about 1 K. While there is at present no generally accepted model of their behavior at higher temperatures, for the past fifteen years the low-temperature behavior has been attributed to the existence in these materials of tunneling two-level systems. The two-level system model is generally believed to receive strong support from a series of remarkable effects (phonon echoes, etc.) observed in the acoustic behavior of these materials.

In this Comment we point out (a) that while the observed effects are (mostly) consistent with the two-level system model, they do not uniquely establish it, (b) that the model cannot explain the dramatic *quantitative* universality in the low-temperature behavior of amorphous materials, and (c) that the neglect to lowest order within the model of the long-range interactions induced by the strain field, while self-consistent, is not imposed by the data. We sketch an alternative scenario for the properties of amorphous materials in which the role played by these interactions is essential. This scenario, if it can be quantitatively implemented, holds out promise of being able to explain in a unified way not only the quantitative universality below 1 K but also the behavior at higher temperatures.

*Comments Cond. Mat. Phys.*  
1988, Vol. 14, No. 4, pp. 231-251

# Comment Paper by Clare and Tony 1988

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*Comments Cond. Mat. Phys.*  
1988, Vol. 14, No. 4, pp. 231-251

# ... a totally unbelievable degree of chance coincidence

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Physica B 169 (1991) 322–327  
North-Holland

## Amorphous materials at low temperatures: why are they so similar?

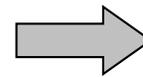
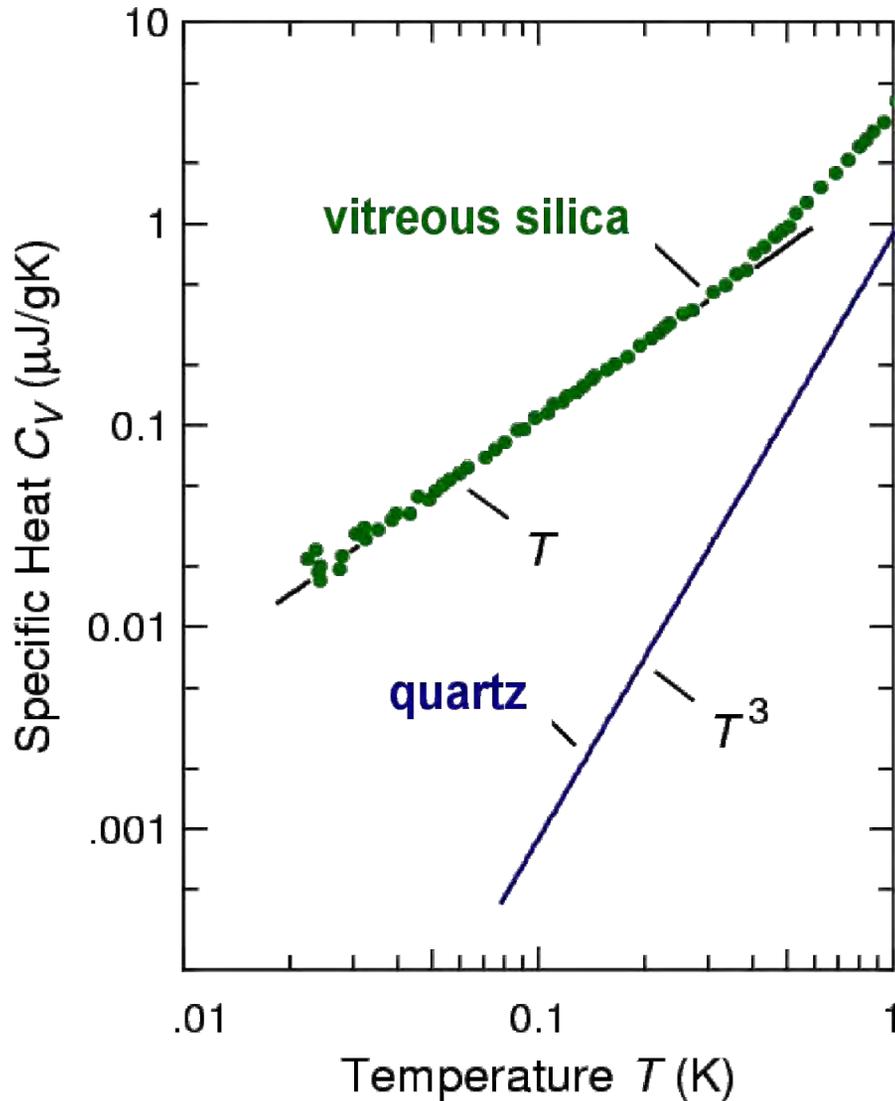
Anthony J. Leggett

*Department of Physics, University of Illinois at Urbana-Champaign, 1110 West Green Street, Urbana, IL 61801, USA*

The invited talk was given by the author.

In addition to the well-known qualitative similarity of behavior of a wide range of amorphous materials, one dimensionless quantity, the reduced attenuation of transverse ultrasound, shows a *quantitative* universality whose explanation in terms of the standard “tunnelling two-level system” model would seem to require a totally unbelievable degree of chance coincidence. In addition, while the height of the “plateau” in the thermal conductivity is material-specific, the higher-temperature behavior is consistent with universal behavior. I sketch the outlines of a scenario which holds out hope of understanding these observations.

# Specific Heat of Amorphous and Crystalline SiO<sub>2</sub>

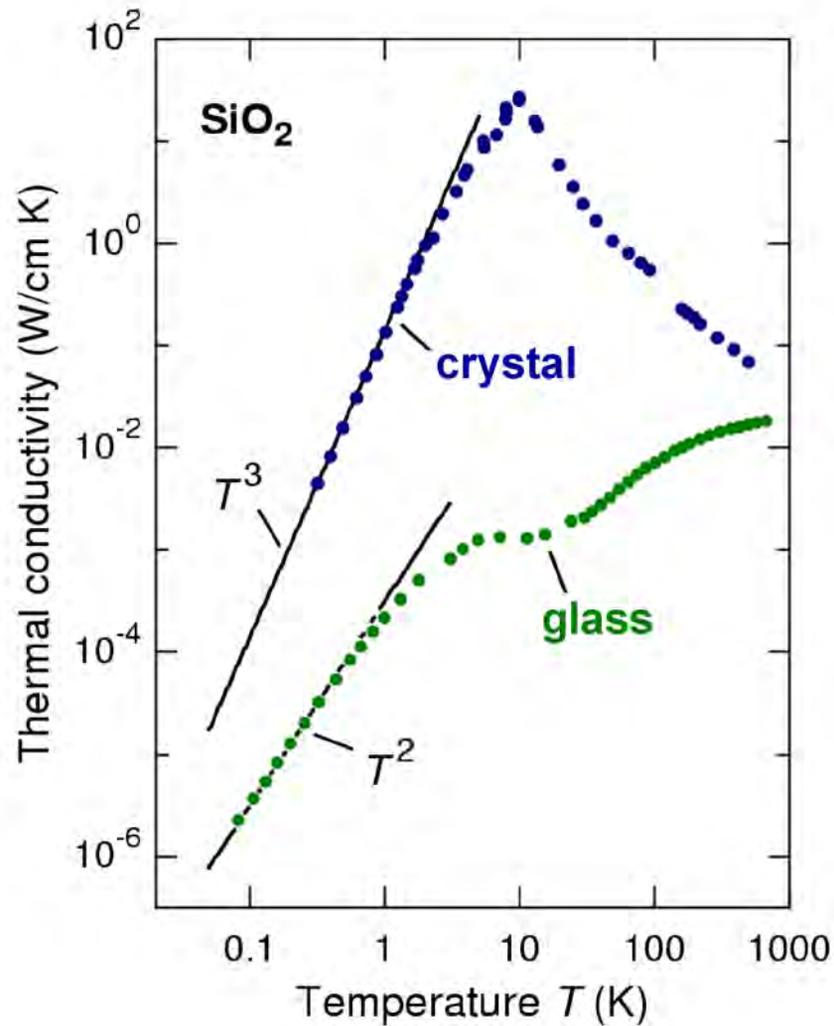


broad distribution of low-energy excitations

R.C. Zeller, R.O. Pohl,  
Phys. Rev. **B 4**, 2029 (1971)

J.C. Lasjaunias et al.,  
Sol. State Commun. **17**, 1045 (1975)

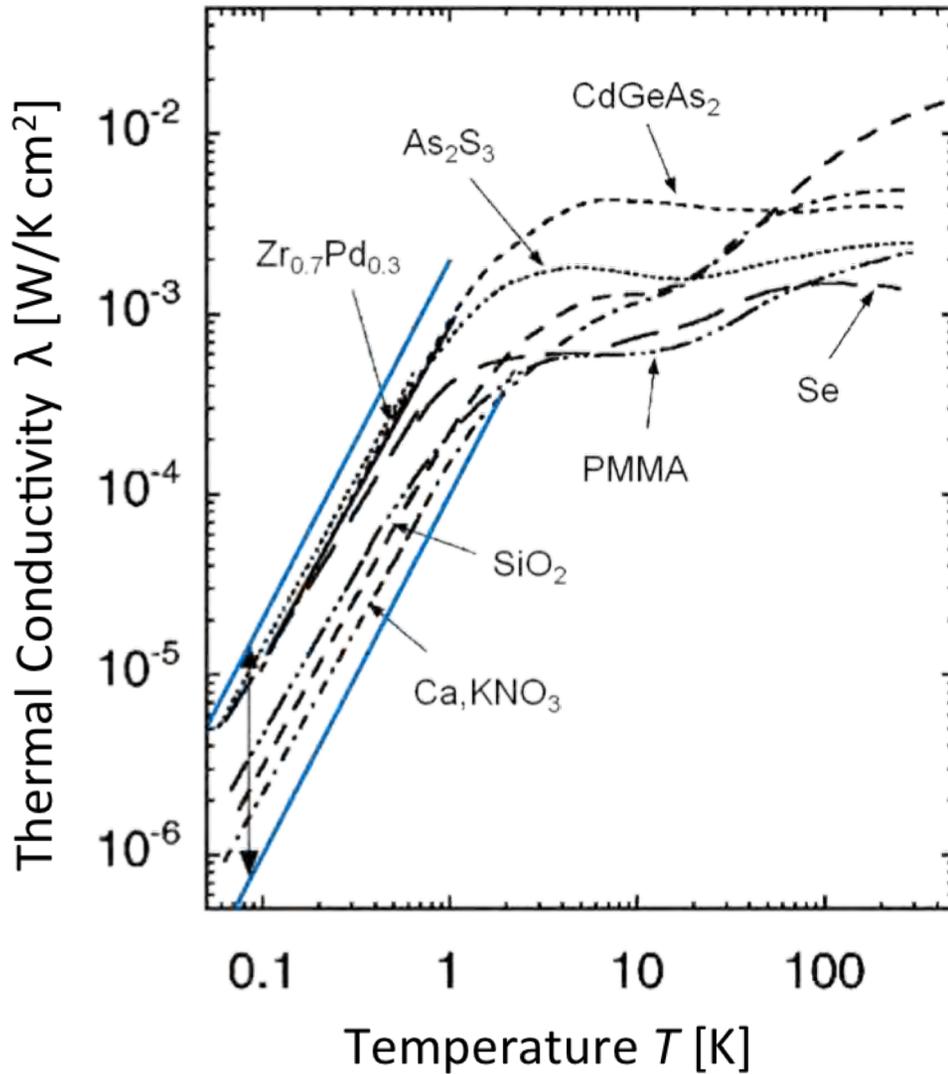
# Specific Heat of Amorphous and Crystalline SiO<sub>2</sub>



- ⇒ excitations are **localized**
- ⇒ **strong coupling** to phonons

R.C. Zeller, R.O. Pohl,  
Phys. Rev. B **4**, 2029 (1971)

# Universality of Glasses



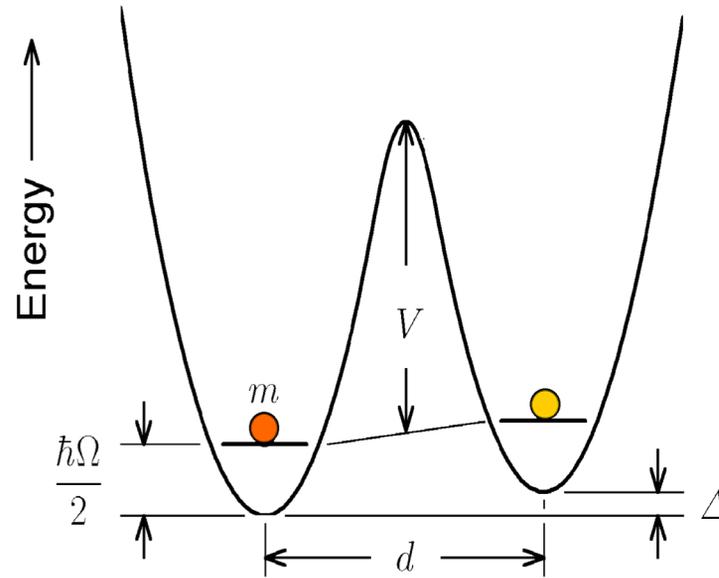
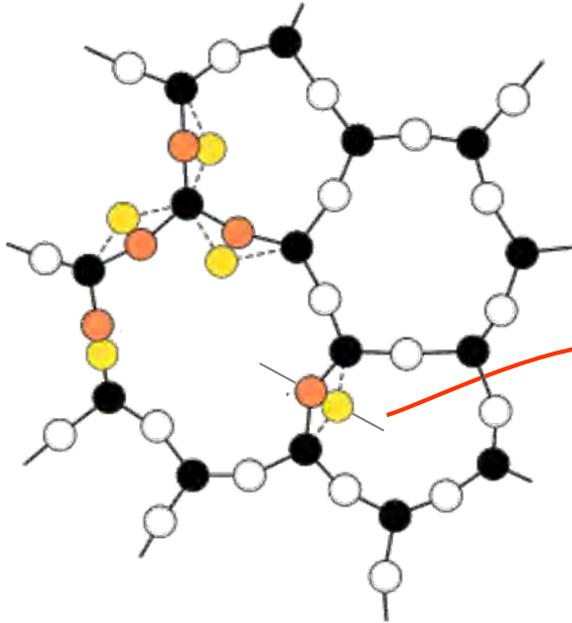
➔ Thermal conductivity of glasses within one order of magnitude

R. B. Stephens, Phys. Rev. B **8**, 2896 (1973)

# Atomic Tunneling Systems in Glasses

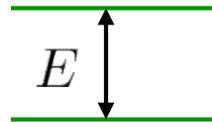
W.A. Phillips, J. Low. Temp. Phys. **7**, 351 (1972)

P.W. Anderson et al., Philos. Mag. **25**, 1 (1972)



energy splitting

$$E = \sqrt{\Delta_0^2 + \Delta^2}$$

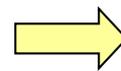


tunnel splitting

$$\Delta_0 = \hbar\Omega e^{-\lambda} \quad \lambda = \frac{d}{2\hbar} \sqrt{2mV}$$

distribution function

$$P(\lambda, \Delta) d\lambda d\Delta = \bar{P} d\lambda d\Delta$$



elastic, dielectric und thermal properities



Dresden 2003



Bibliothek

Armin Fuchs

# Is there an unambiguous test for TLS in glasses?

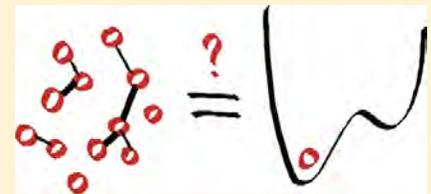
## “Tunneling Two-Level Systems” Model of the Low-Temperature Properties of Glasses: Are “Smoking-Gun” Tests Possible?

Anthony J. Leggett<sup>\*,†</sup> and Dervis C. Vural<sup>\*,‡</sup>

<sup>†</sup>Department of Physics, University of Illinois at Urbana–Champaign, 1110 West Green Street, Urbana, Illinois, United States

<sup>‡</sup>Applied Physics, SEAS, Harvard University, 29 Oxford Street, Cambridge, Massachusetts, United States

**ABSTRACT:** Following a brief review of the “two-level (tunneling) systems” model of the low-temperature properties of amorphous solids (“glasses”), we ask whether it is in fact the unique explanation of these properties as is usually assumed, concluding that this is not necessarily the case. We point out that (a) one specific form of the model is already experimentally refuted and (b) that a definitive test of the model in its most general form, while not yet carried out, would appear to be now experimentally feasible.



*J. Phys. Chem. B* 2013, 117, 12966–12971

# The Cindarella Problem of Glasses

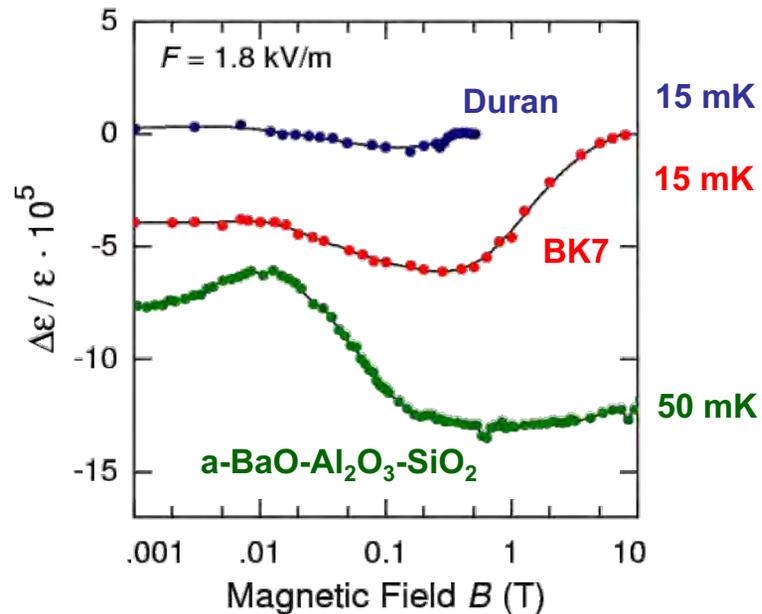
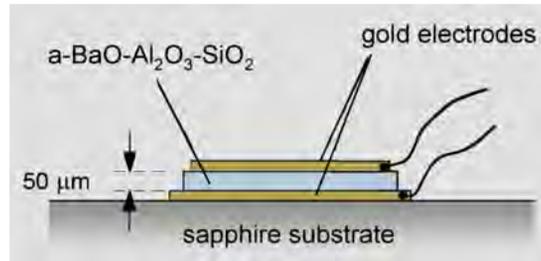
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Universality of glasses cannot be explained with the Standard TLS Model (although it is itself no contradiction to it).

At the same time universality is the very reason, why it is so hard to determine the microscopic nature of the excitations in glasses.

# Unexpected Magnetic Field Dependence

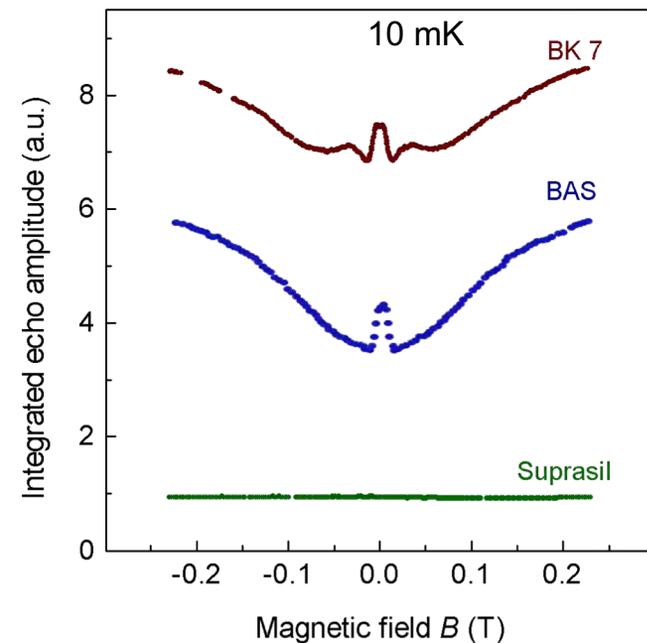
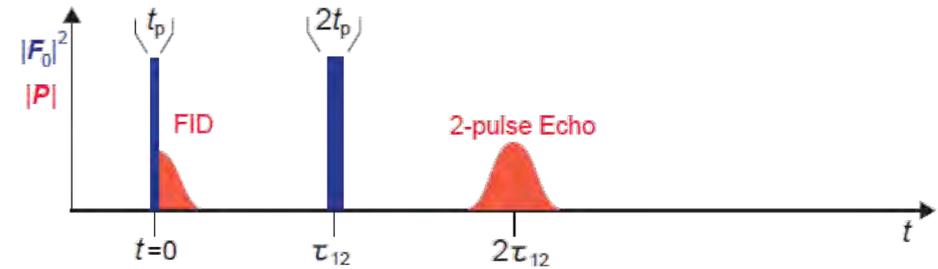
## Dielectric Susceptibility



BK7 and Duran: *M. Wohlfahrt, P. Strehlow, C. E., S. Hunklinger, Europhys. Lett. 56, 690 (2001)*

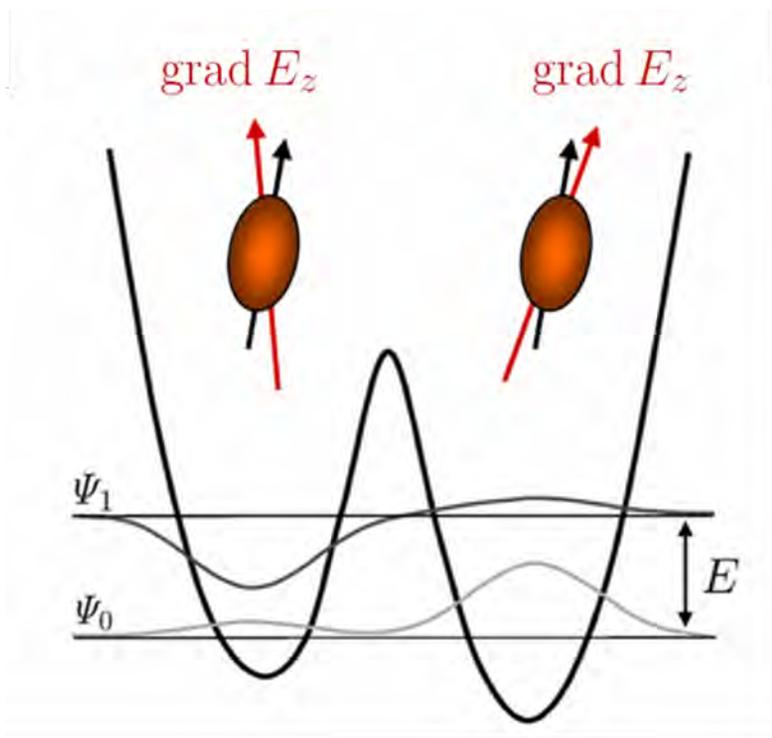
a-BaO- Al<sub>2</sub>O<sub>3</sub>- SiO<sub>2</sub>: *R. Haueisen, G. Weiss, Physica B 316-317, 555 (2002)*

## Dipole Echoes

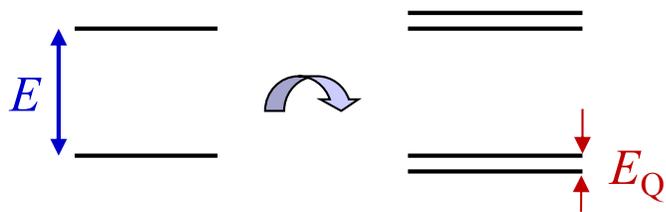


*S. Ludwig, C. E., S. Hunklinger, P. Strehlow, Phys. Rev. Lett. 88, 75501 (2002)*

# Nuclear Quadrupole Moment is Important



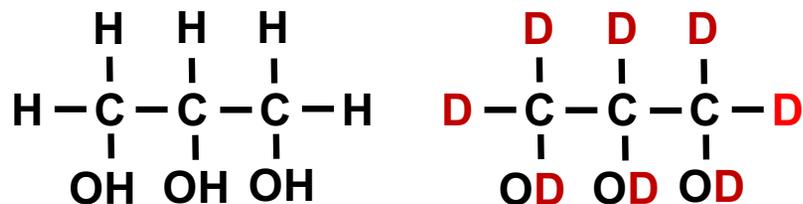
- nuclear quadrupole moment sees electric field gradient in the two wells
  - ➔ hyperfine splitting of tunneling levels
  - ➔ multi-level systems
- magnetic field causes an additional Zeeman splitting of nuclear levels
- no effect for  $\alpha\text{-SiO}_2$  because no quadrupole moment



A. Würger, A. Fleischmann, C. E.,  
Phys. Rev. Lett. **89**, 237601 (2002)

# Proof: Isotope Effect $H \leftrightarrow D$

Glycerol



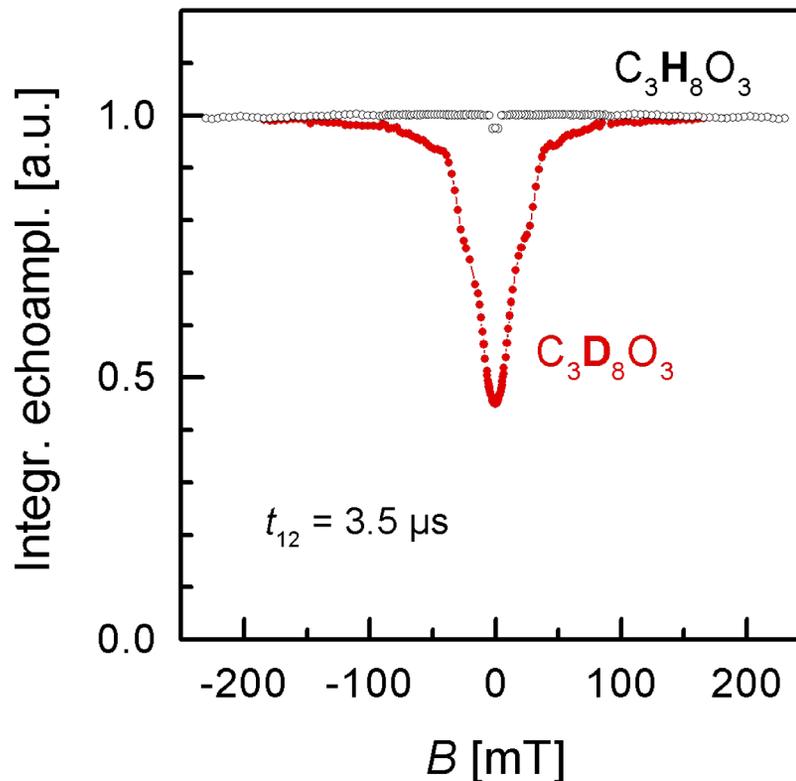
hydrogen

$$I = 1/2, \quad \mu = 2.79 \mu_N, \quad Q = 0$$

deuterium atom

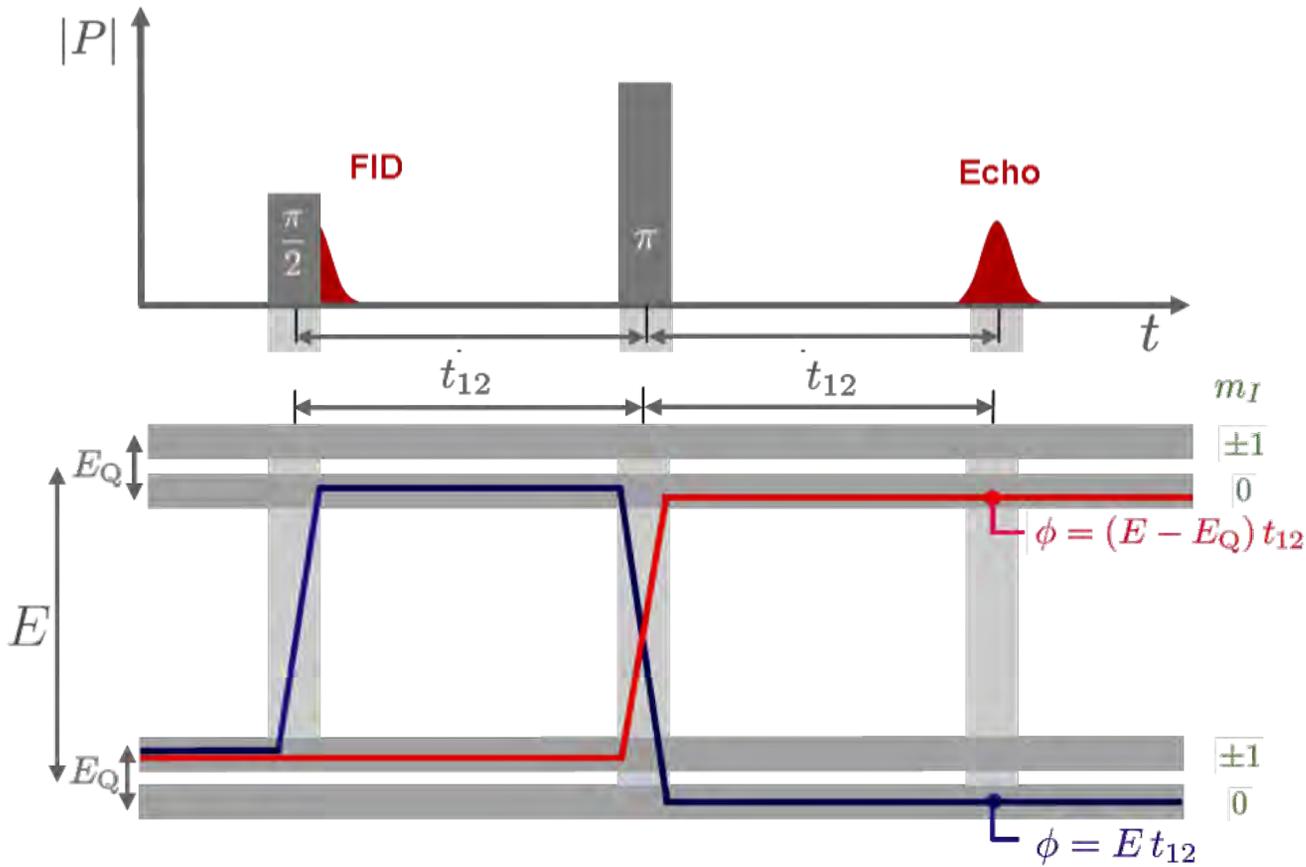
$$I = 1, \quad \mu = 0.86 \mu_N, \quad Q = 0.0029 b$$

→ proof of the quadrupole model



P. Nagel, A. Fleischmann, S. Hunklinger, C. E.,  
Phys. Rev. Lett. **92**, 24511-1 (2004)

# Zero Magnetic Field



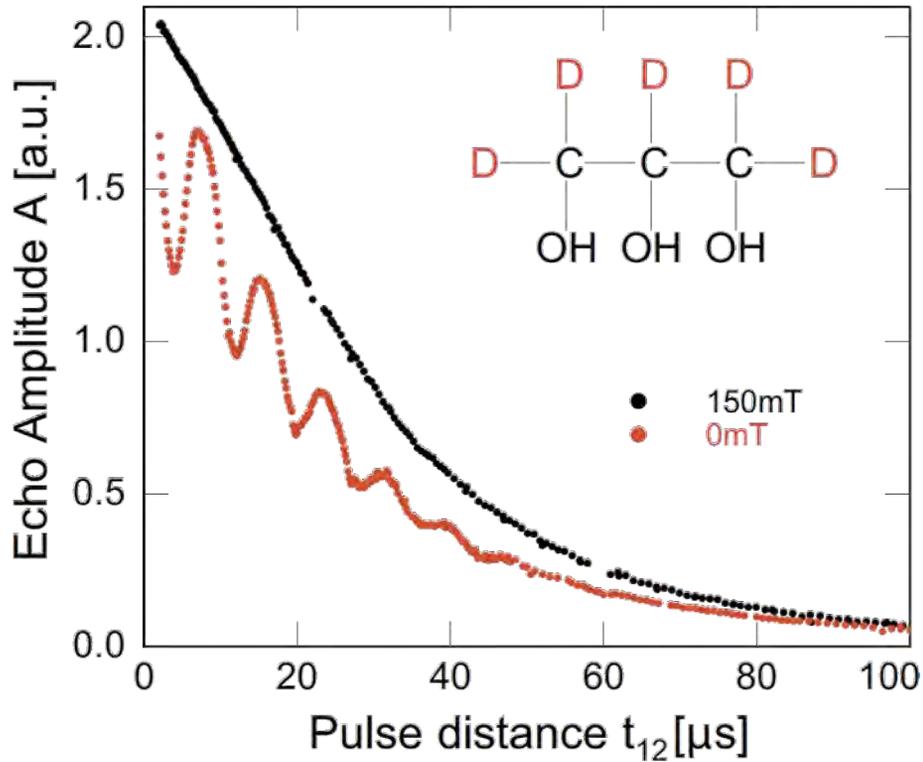
A. Würger, JLTP 137, 143 (2004)

D.A. Parshin, JLTP 137, 233 (2004)

$$A = A_0 [a_1 + a_2 \cos(\omega_Q t_{12}) + a_3 \cos(2\omega_Q t_{12})]$$

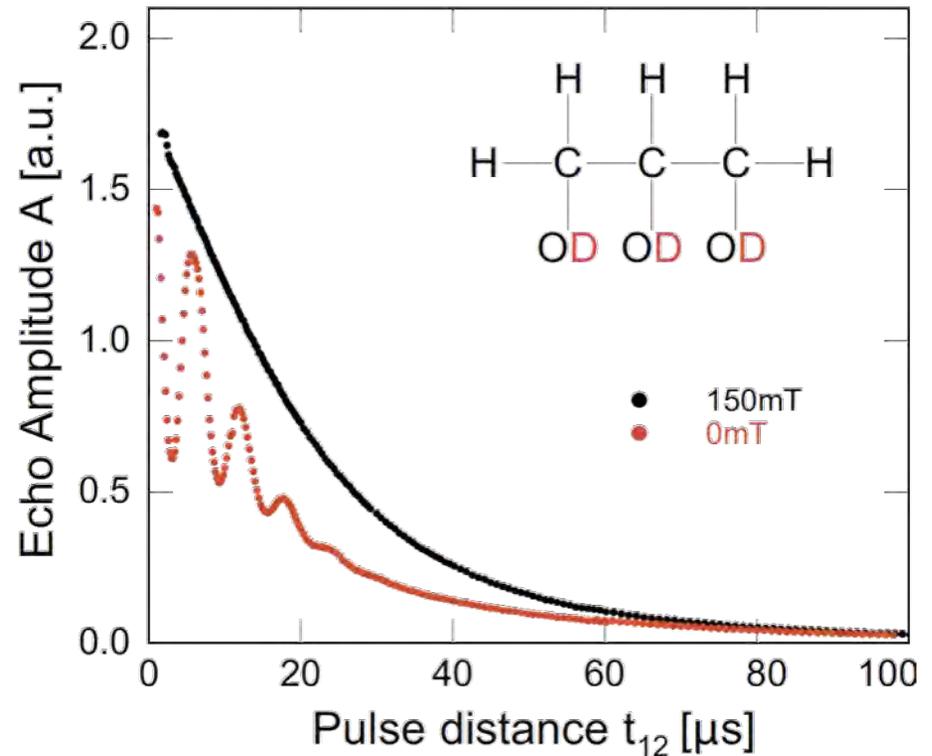
$$a_1 + a_2 + a_3 = 1$$

# Partially Deuterated Glycerol



$$\nu_Q = 128 \text{ kHz}$$

$$\nu_Q = 125 \text{ kHz (NMR)}$$



$$\nu_Q = 160 \text{ kHz}$$

$$\nu_Q = 158 \text{ kHz (NMR)} \quad \text{W. Schnauss, F. Fujara, H. Sillescu, J. Chem. Phys. 97, 1378 (1992).}$$

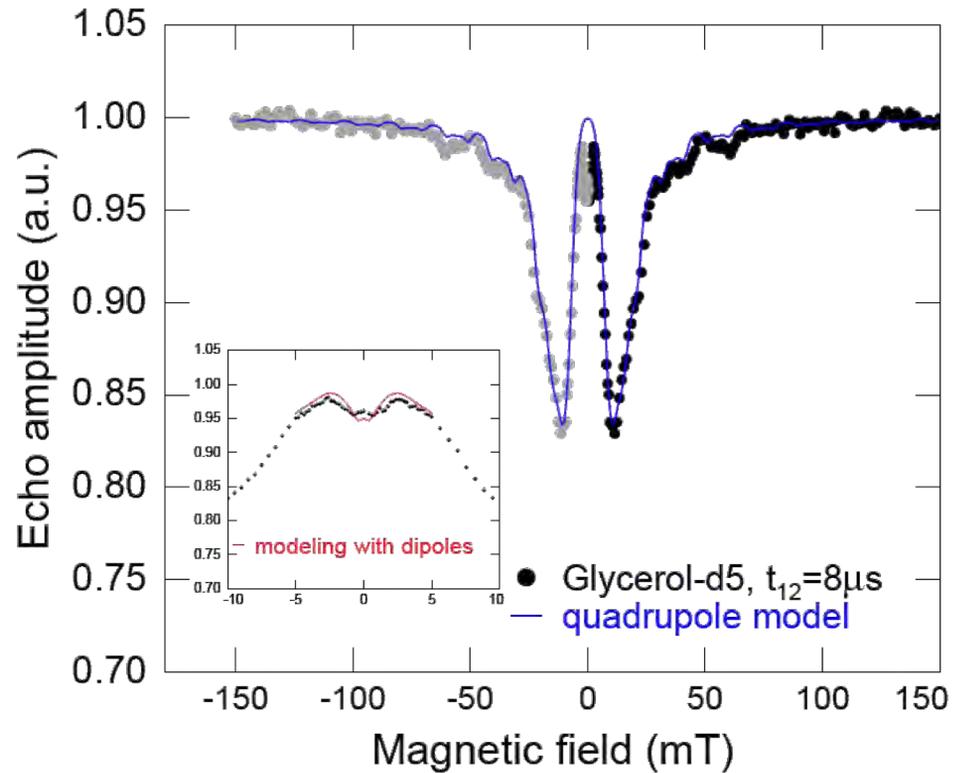
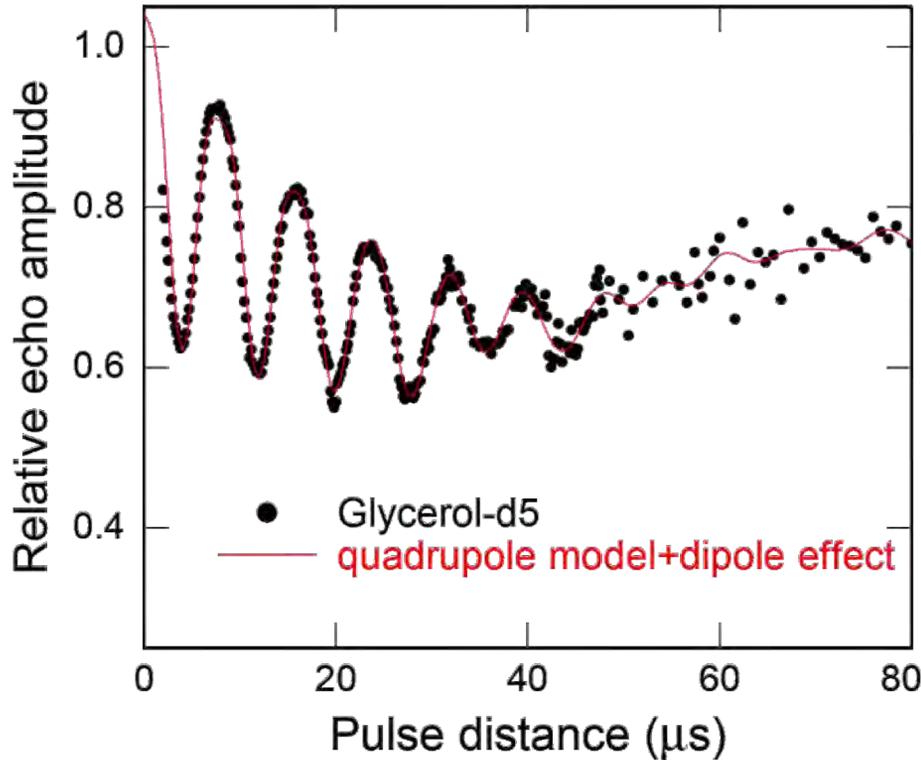
Beating in case of  $d_3$  disappears faster



local environment ?

A. Bartkowiak, M. Bazrafshan, C. Fischer, A. Fleischmann, C. E., Phys. Rev. Lett. 110 (2013)

# Microscopic Modeling



Tunneling angle  $2\theta = 17^\circ$

M. Bazrafshan, PhD Thesis 2008

A. Bartkowiak, M. Bazrafshan, C. Fischer, A. Fleischmann, C. E., Phys. Rev. Lett. **110** (2013)



First measurement of **microscopic** properties of TLS in Glass

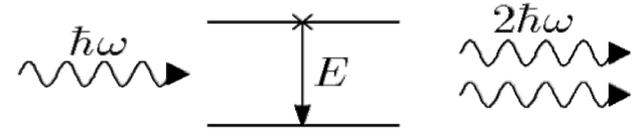
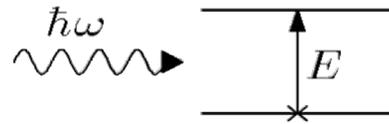
# Dielectric Susceptibility

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Can nuclear spins explain the magnetic field dependence of the dielectric susceptibility?

# Coupling to Electric and Elastic Fields

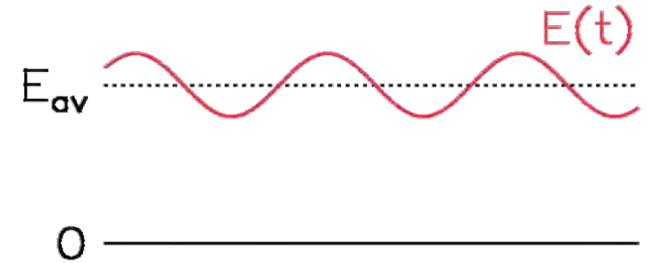
resonant processes



relaxational processes

→ modulation of  $\Delta$       $\delta\Delta = 2\gamma e$

$$\delta\Delta = 2\mathbf{p}\cdot\mathbf{F}$$

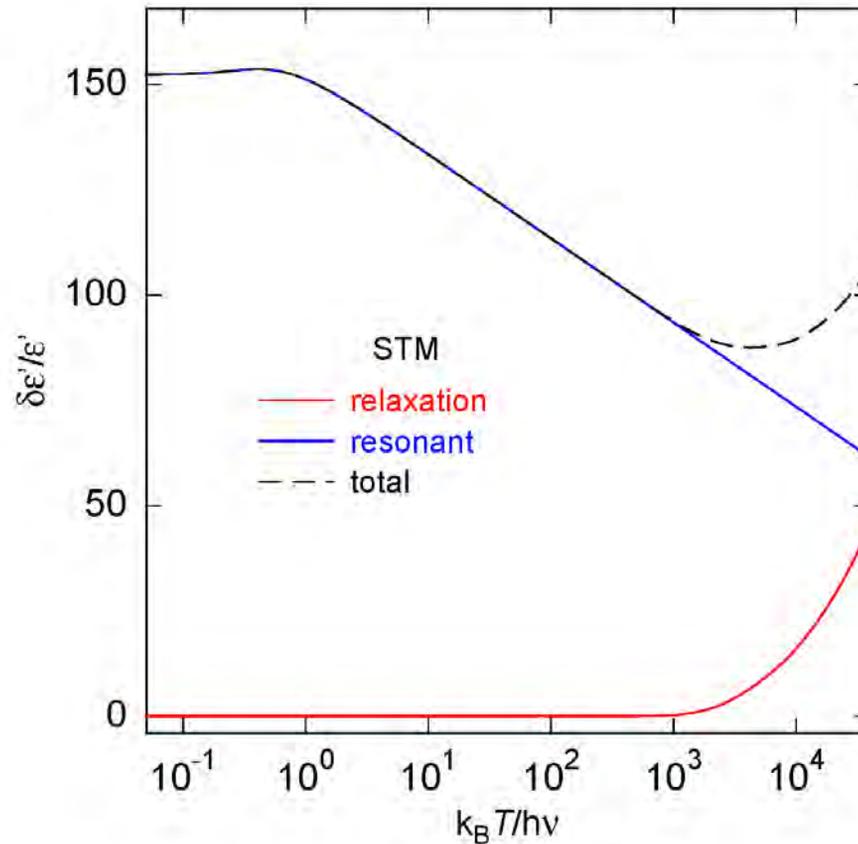


$T < 1$  K one-phonon relaxation

$$\tau_1 = A \left( \frac{E}{\Delta_0} \right)^2 \frac{1}{E^3} \tanh \left( \frac{E}{2k_B T} \right)$$

# Dielectric Susceptibility in the STM

Real part after integration of tunneling parameter distribution

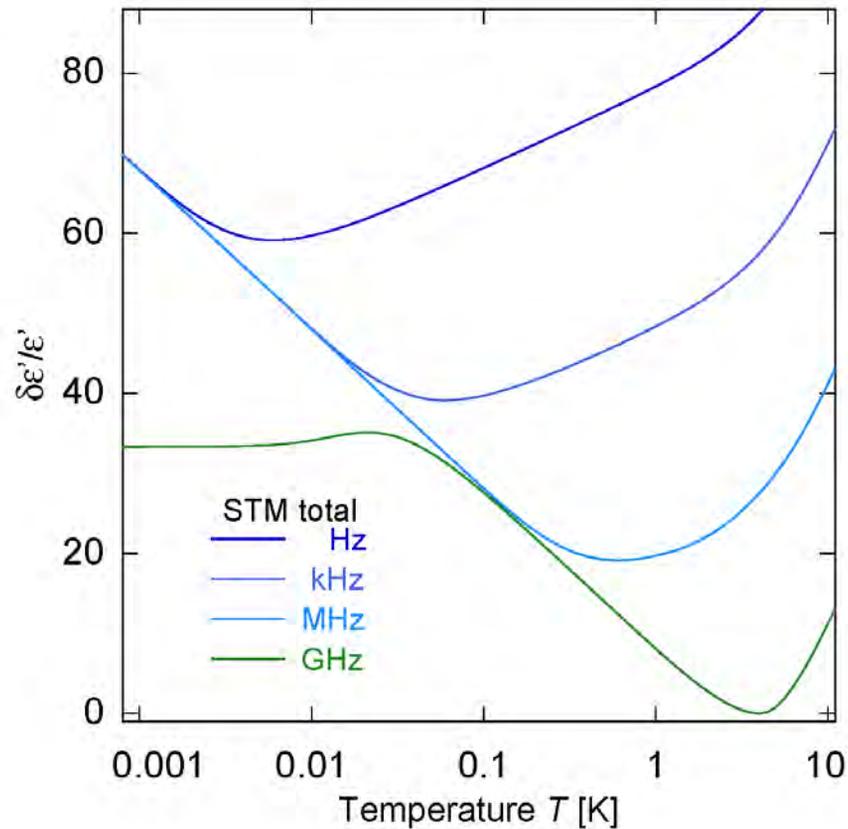


# Different Frequencies

## Relaxation Rate

$$\tau_{\text{tot}}^{-1} = \tau_{1\text{ph}}^{-1} + \tau_{2\text{ph}}^{-1}$$

$\tau_{1\text{ph}}^{-1} \propto E^3 \propto T^3 \longrightarrow T_{\text{min}} \propto f^{1/3}$  one phonon  
 $\tau_{2\text{ph}}^{-1} \propto T^7 \longrightarrow T_{\text{min}} \propto f^{1/7}$  two phonon

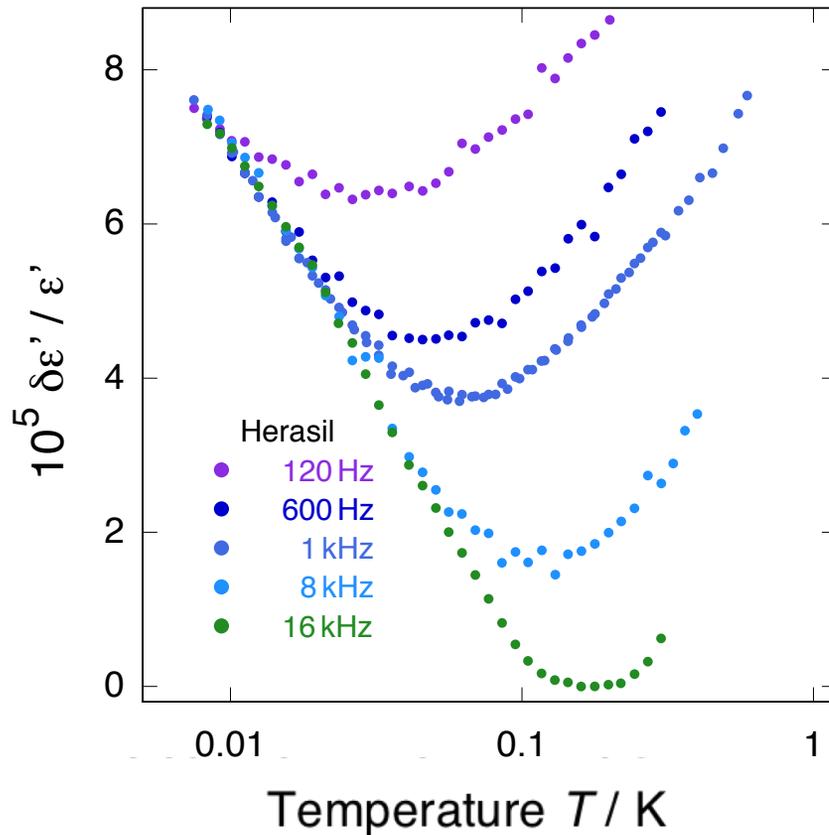


# Samples

Name	Type of Glass	Quadrupole Moment $Q$
CVAc	Polymer glass	None
Herasil	Pure quartz glass	None
BK7	Multi-component glass	No large $Q$
N-KZFS11	Multi-component glass	$^{181}\text{Ta}$ , $Q=3.3$ barn
HY-1	Multi-component glass	$^{165}\text{Ho}$ , $Q=3.5$ barn

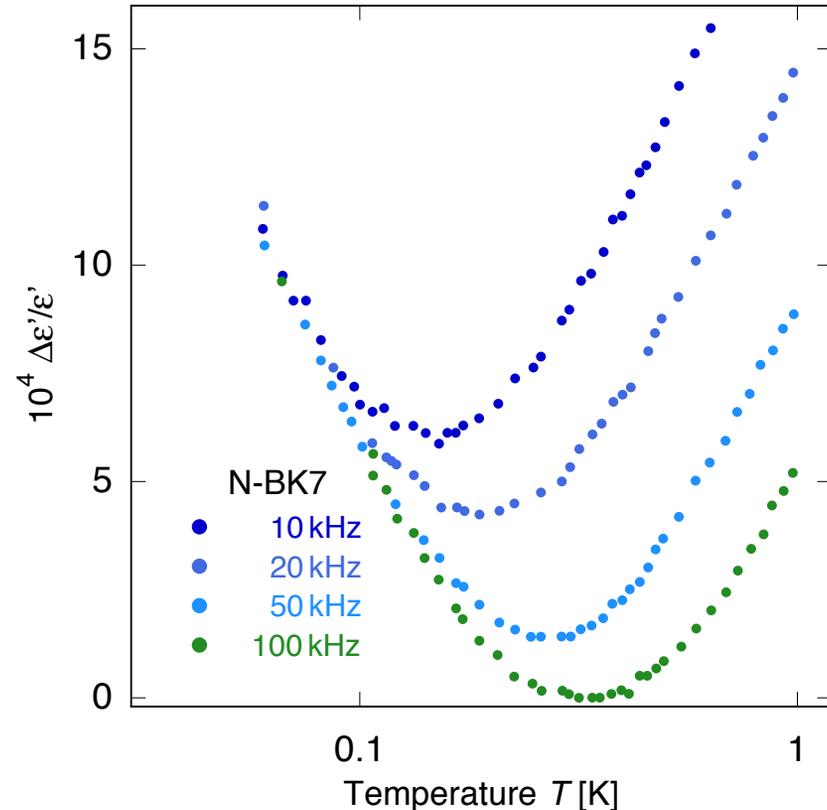
# Dielectric Susceptibility of Standard Glasses

Herasil



A. Luck, A. Fleischmann, A. Reiser, C. E.,  
J. Phys. C **568**, 032013 (2014)

Multi-component glass BK7



C.E., C. Bechinger, M. v. Schickfus,  
Phonons **89**, 474 (1989)

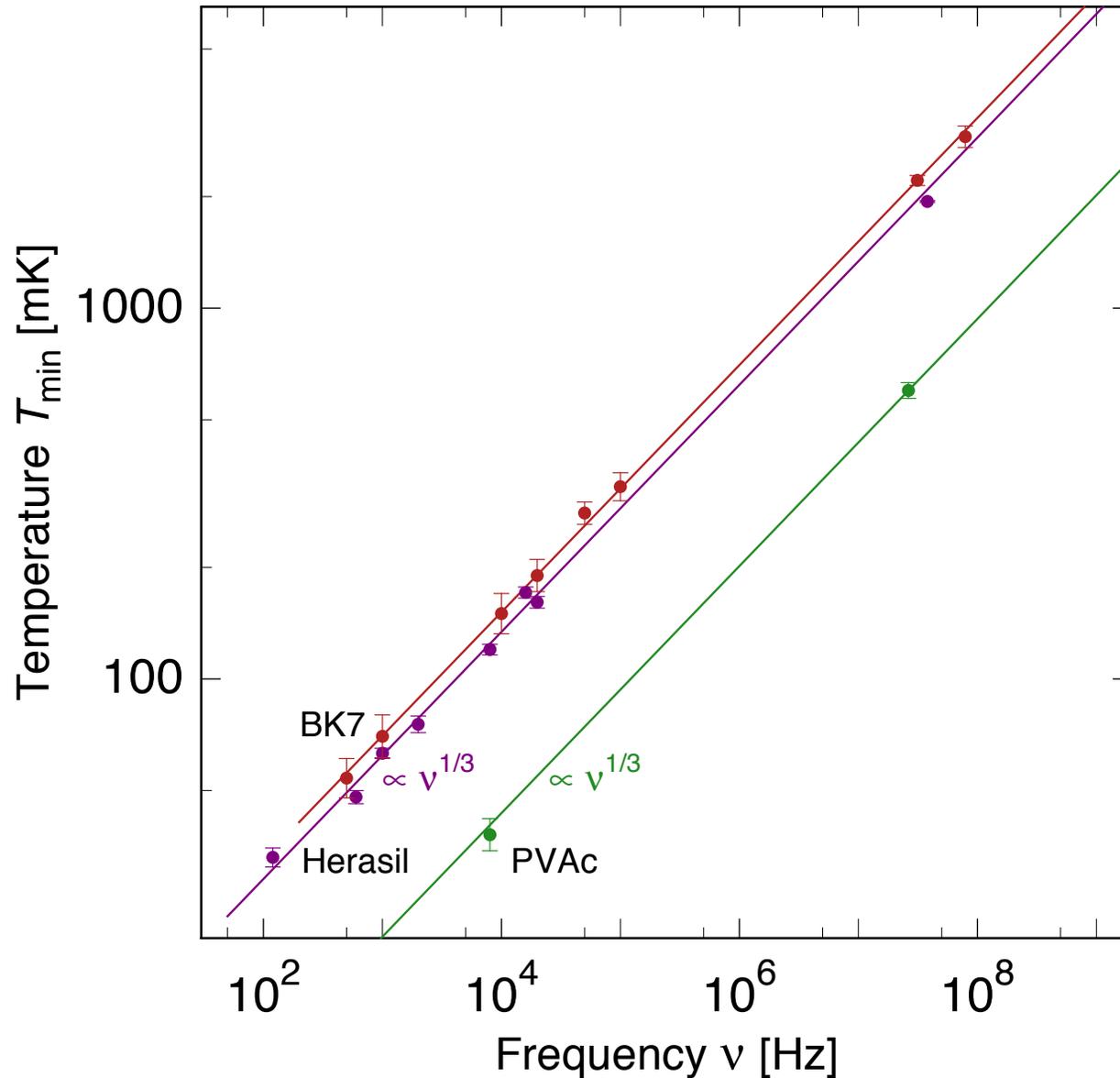


Logarithmic temperature dependence



shift of minimum temperature with frequency

# Frequency Dependence of Minimum in DK



one phonon process:

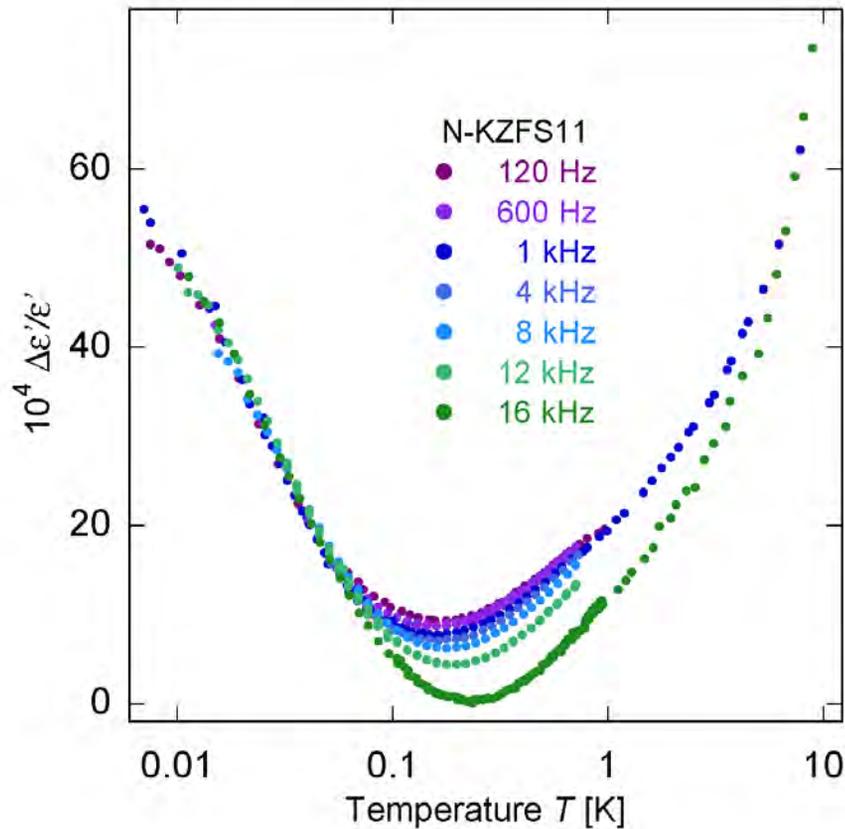
$$T_{\min} \propto \nu^{1/3}$$

two phonon process:

$$T_{\min} \propto \nu^{1/7}$$

# Dielectric Susceptibility of Glasses With Nuclear Quadrupols

Multi-component glass containing Tantalum

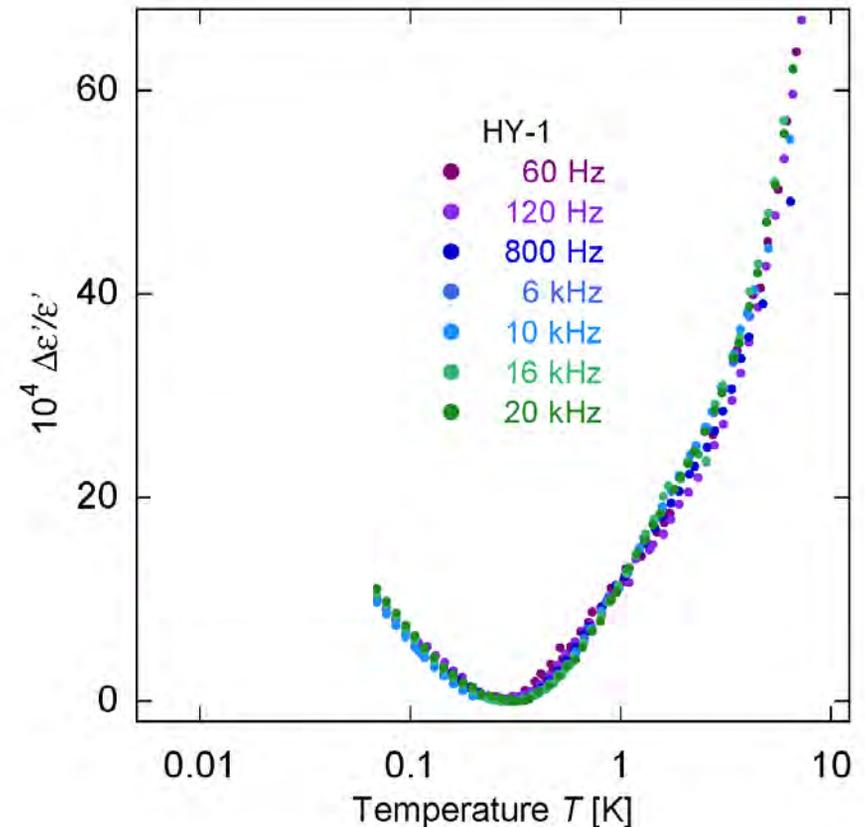
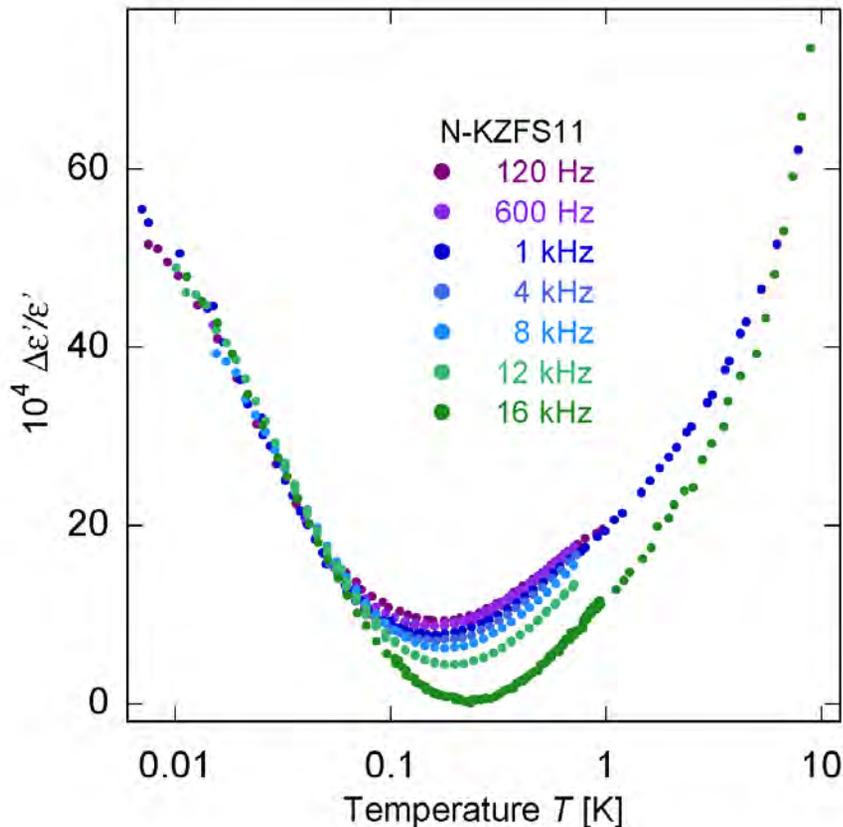


A. Luck, A. Fleischmann, A. Reiser, C. Enss,  
J. Phys. C **568**, 032013 (2014)

# Dielectric Susceptibility of Glasses With Nuclear Quadrupols

Multi-component glass containing **Tantalum**

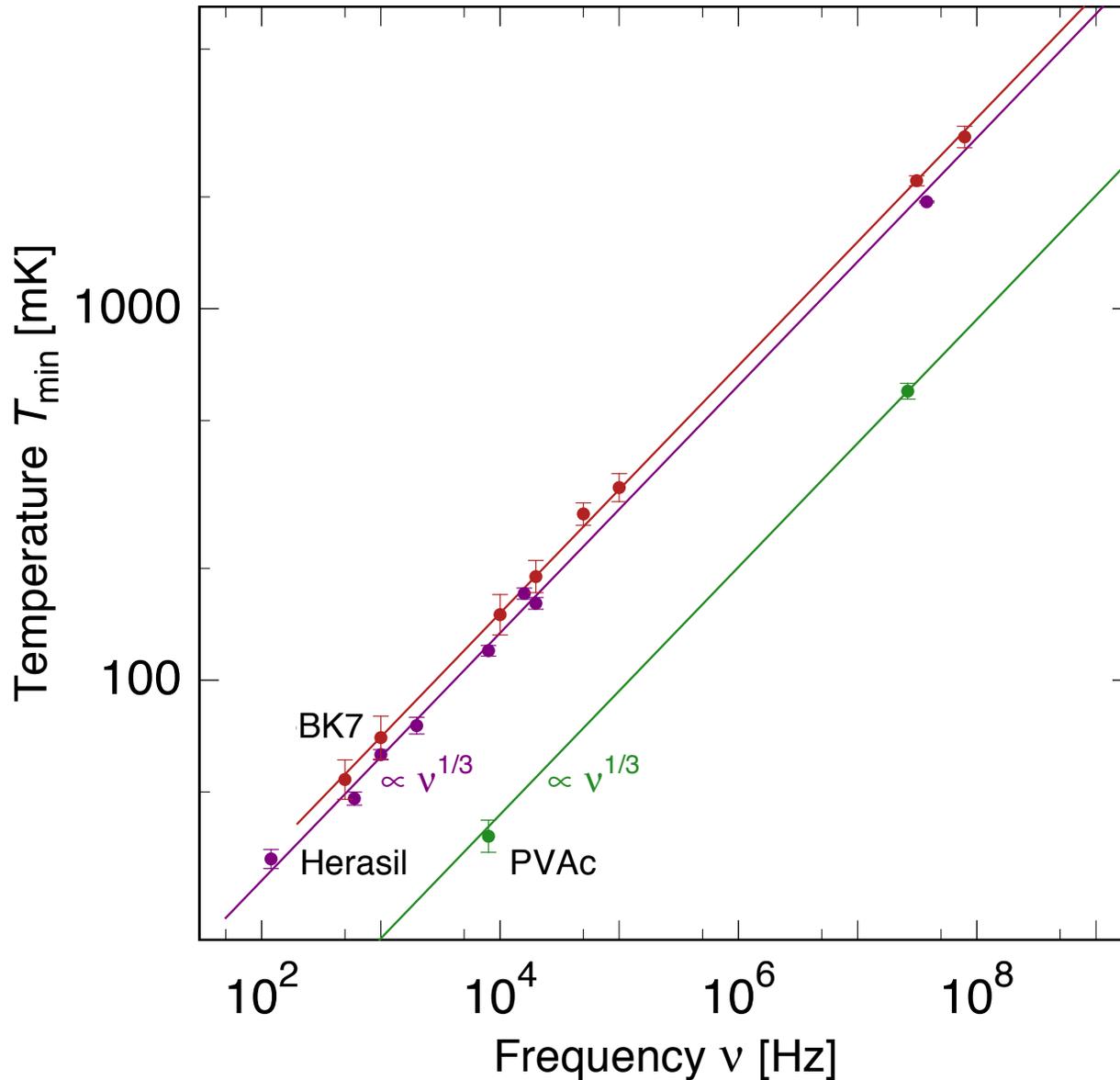
Multi-component glass containing **Holmium**



A. Luck, A. Fleischmann, A. Reiser, C. Enss, *J. Phys. C* **568**, 032013 (2014)

- ➡ HY-1 and N-KZFS11 **different** from other glasses
- ➡ no shift of **minimum temperature** with **frequency** at low frequencies

# Frequency Dependence of Minimum in DK



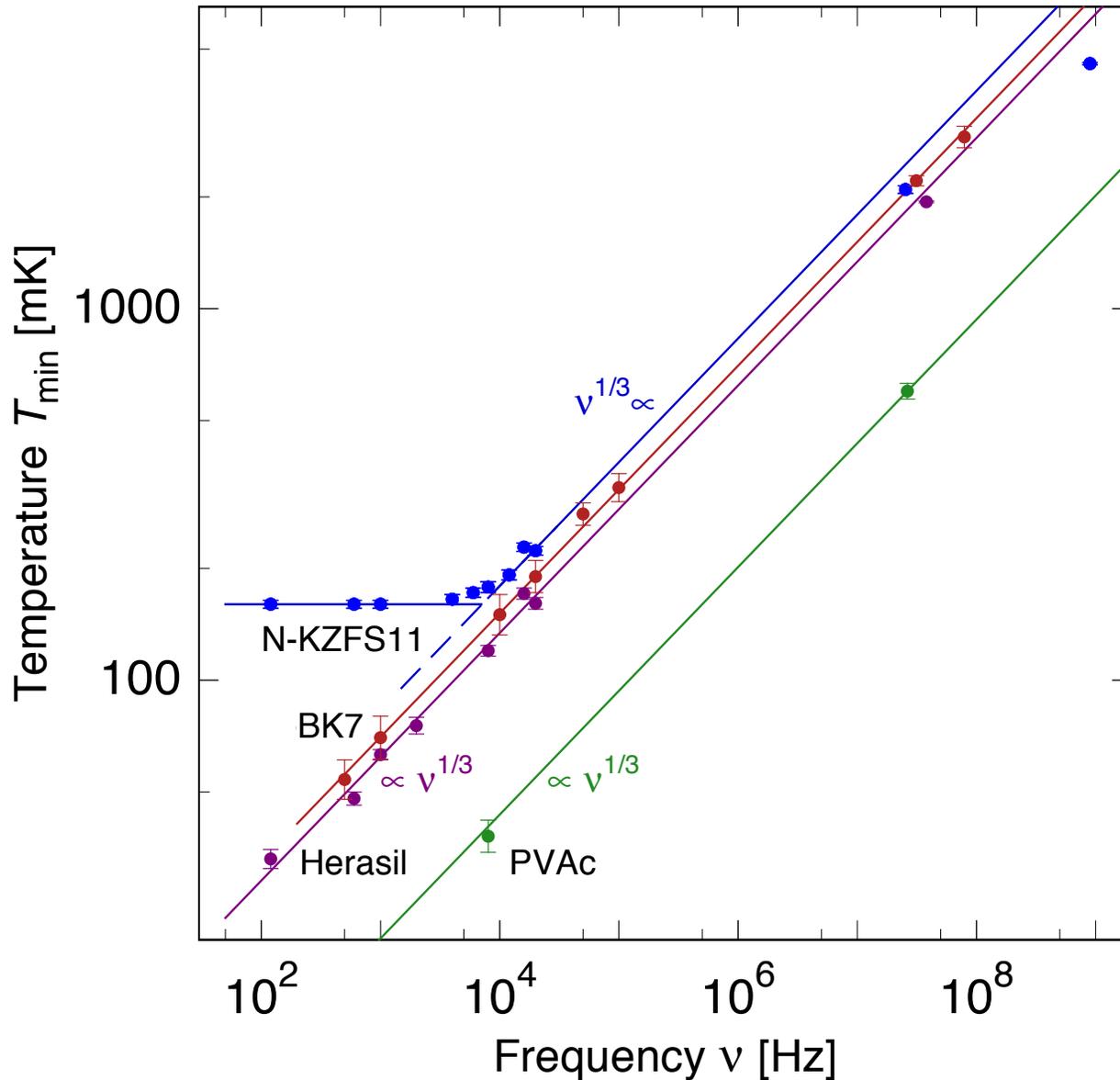
one phonon process:

$$T_{\min} \propto \nu^{1/3}$$

two phonon process:

$$T_{\min} \propto \nu^{1/7}$$

# Frequency Dependence of Minimum in DK



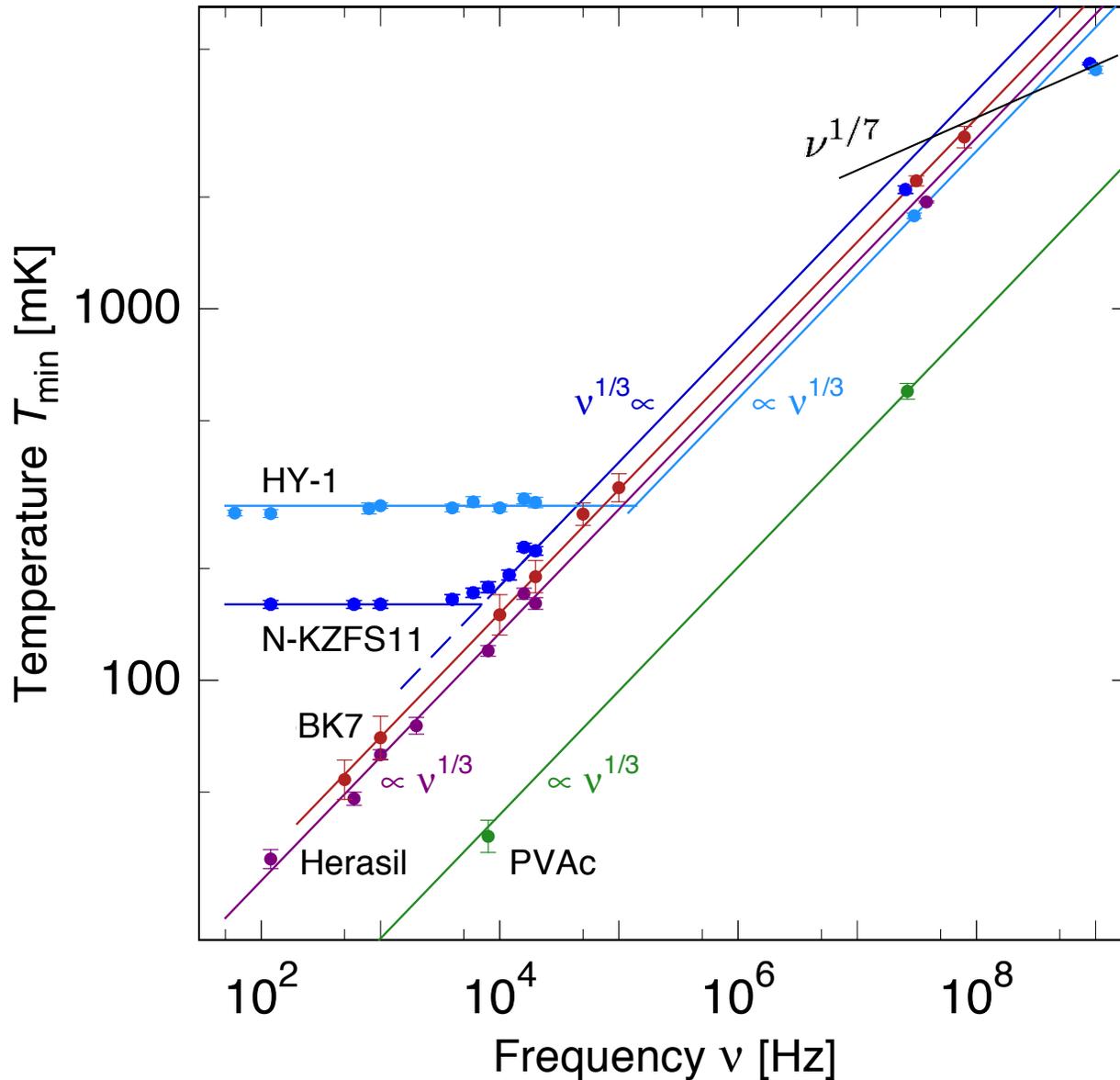
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one phonon process:

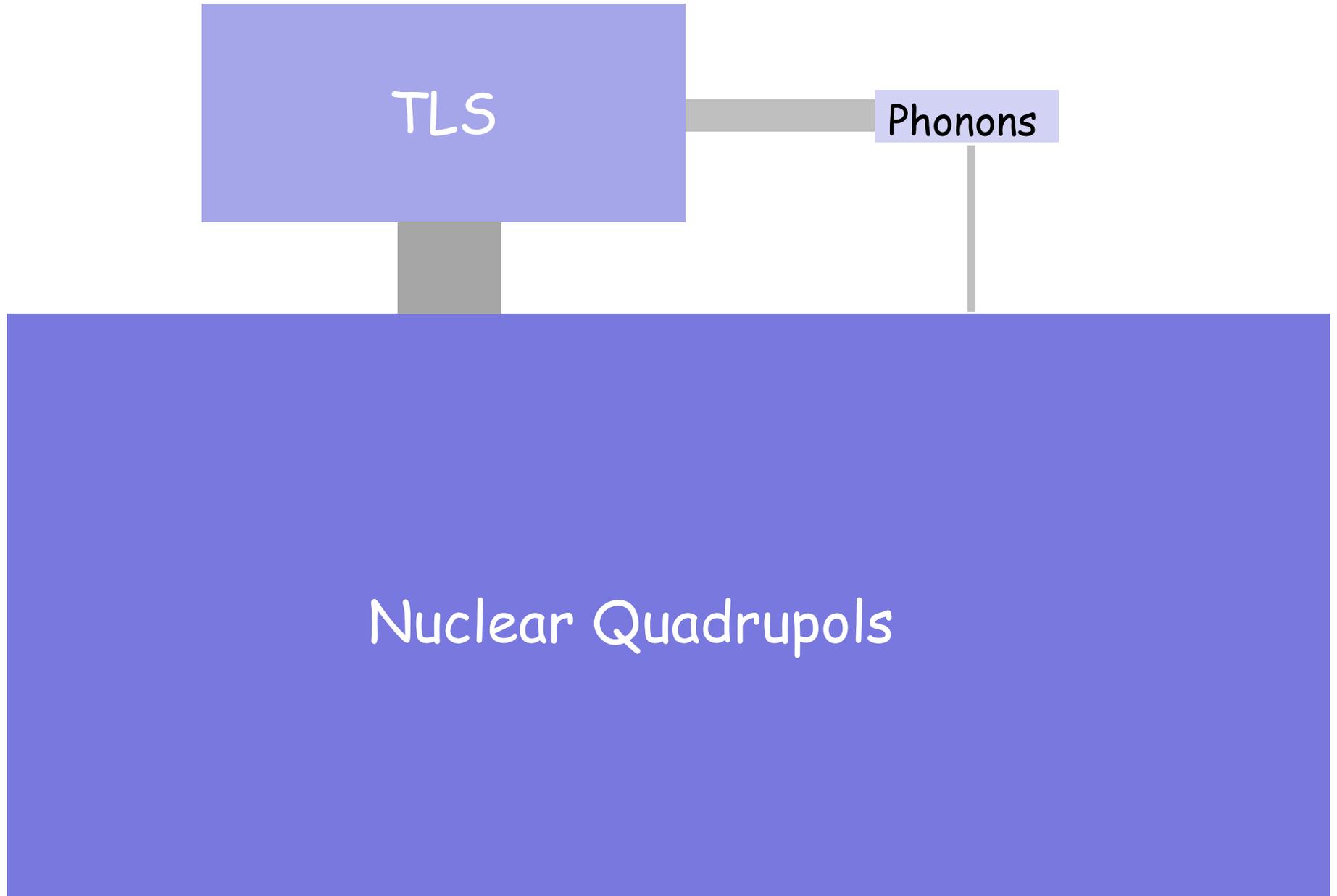
$$T_{\min} \propto \nu^{1/3}$$

two phonon process:

$$T_{\min} \propto \nu^{1/7}$$

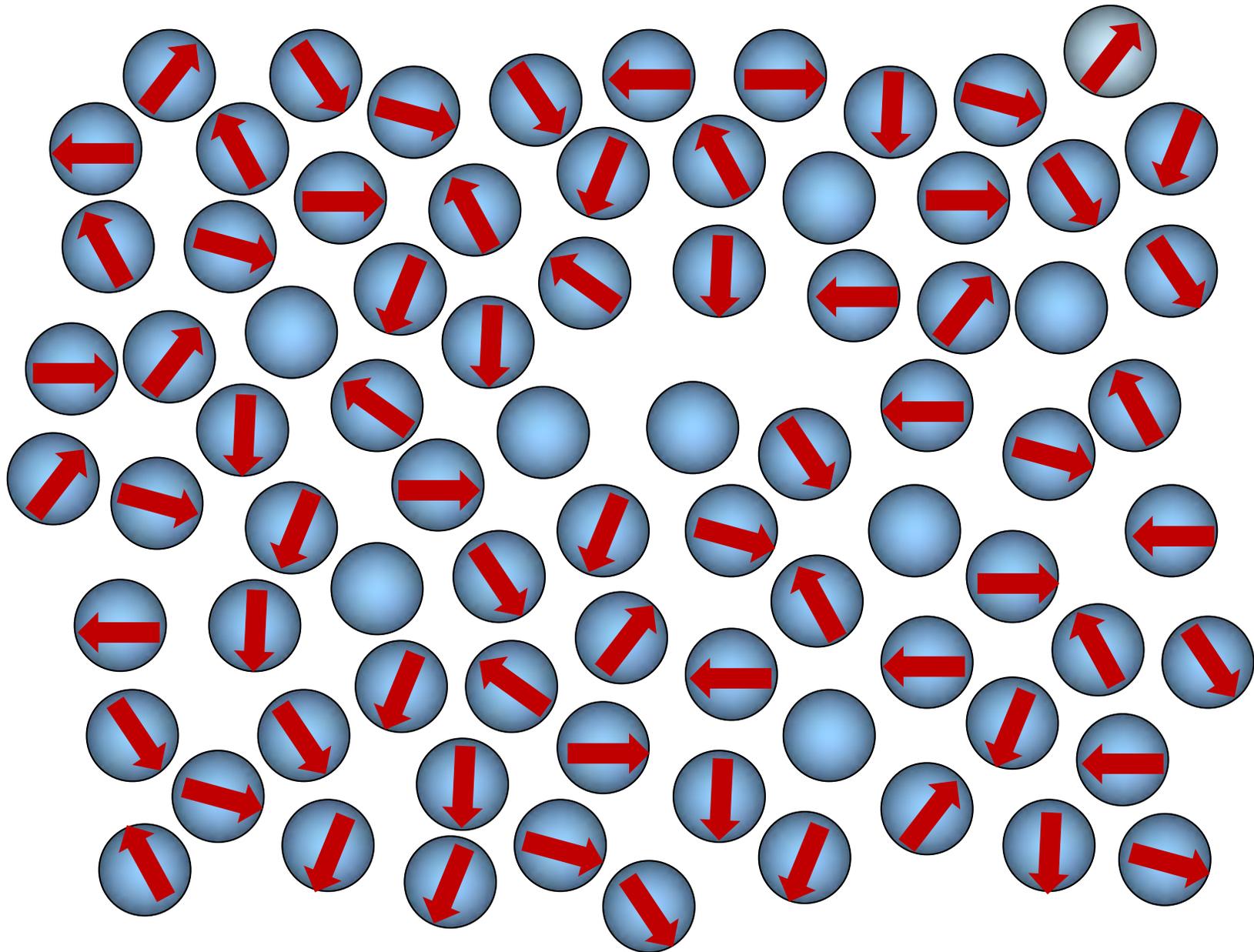
# Relaxation Into Nuclear Spin Bath

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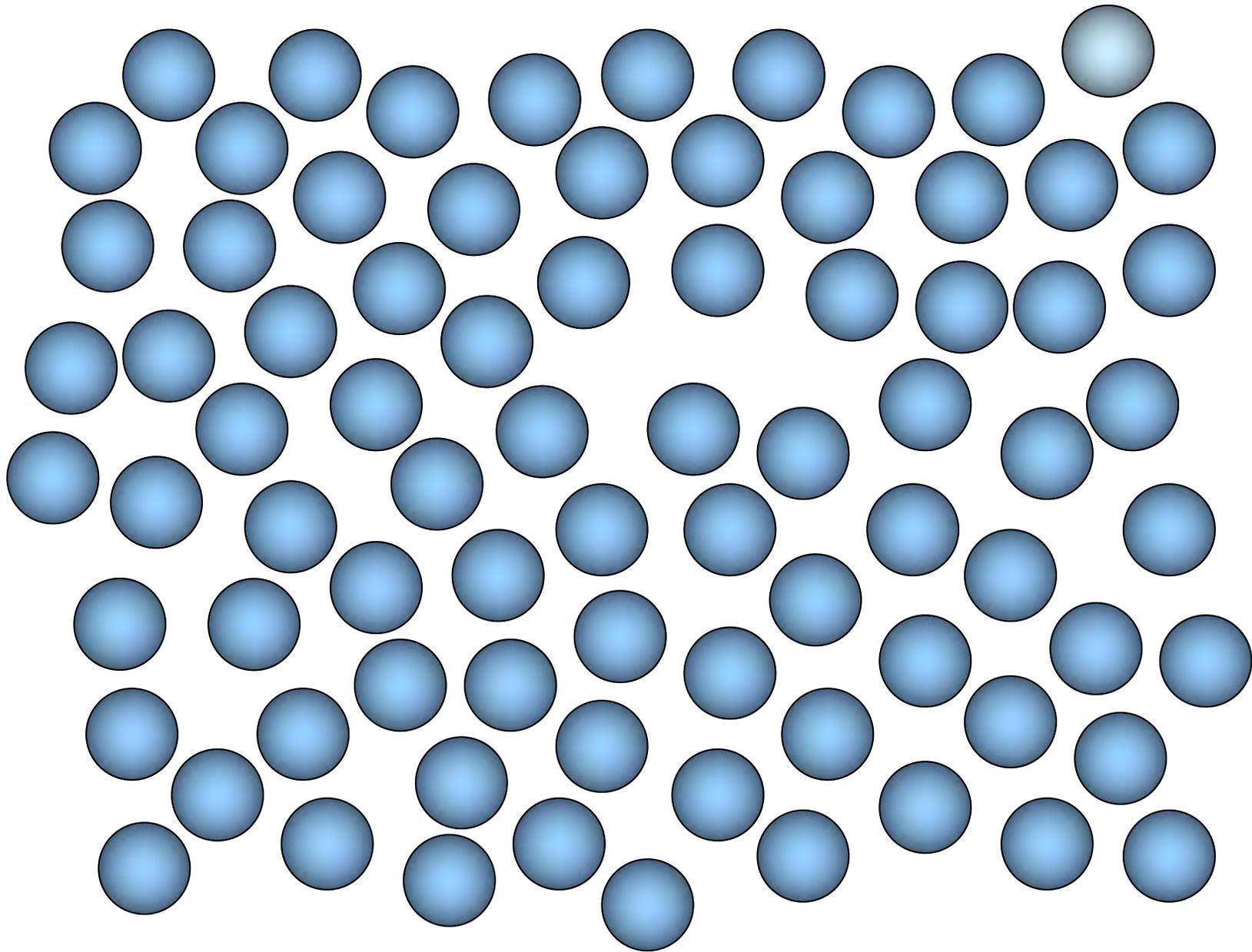
# Relaxation Into Nuclear Spin Bath

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# Atomic Tunneling Systems in Glasses

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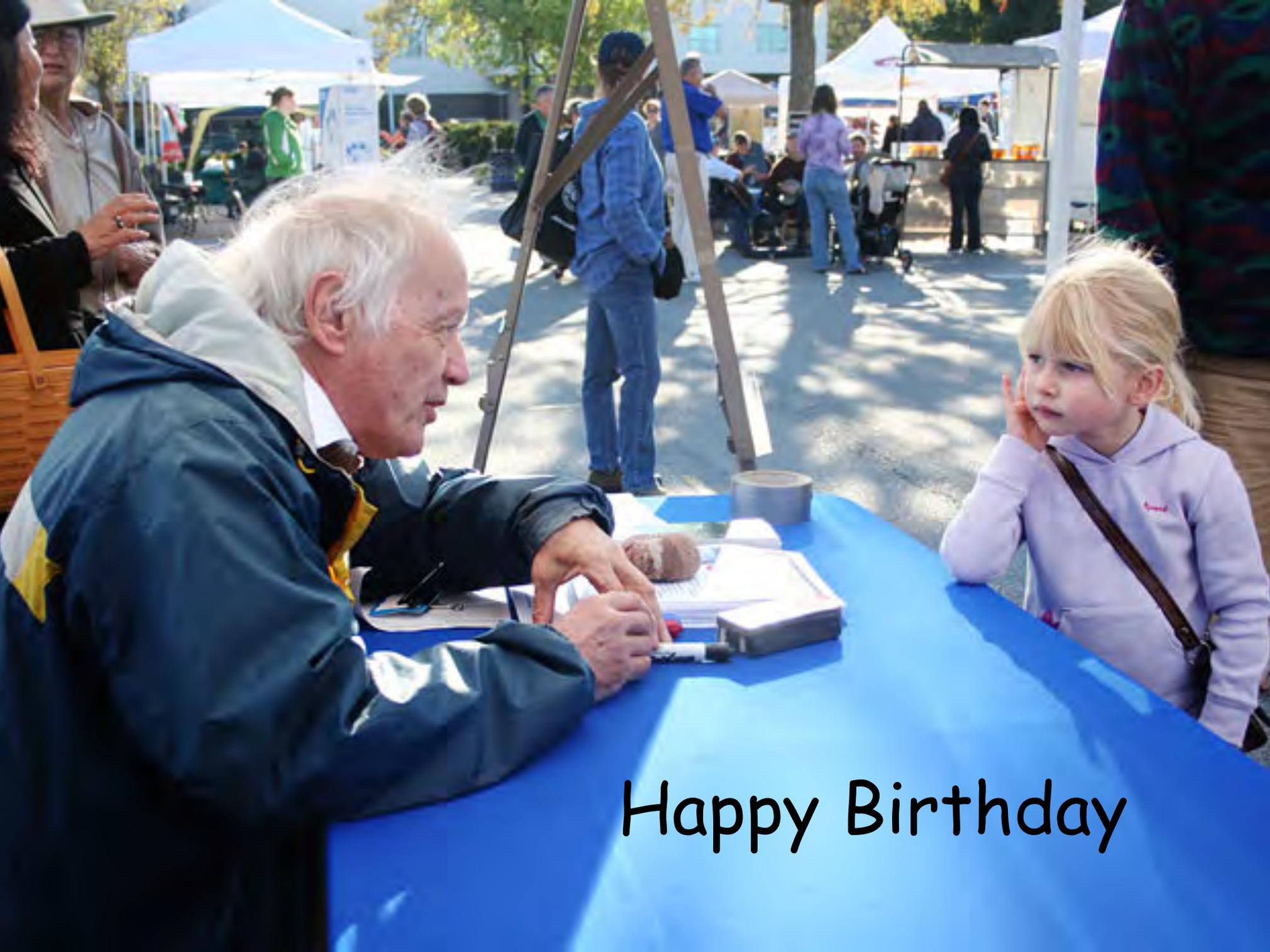


# Summary

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- Nuclear spins are important in glasses at low temperatures
- The universality of glasses breaks down at very low temperatures
- This allows for material dependent studies
- TLS model can be underpinned microscopically at least for some materials
- Still no microscopic explanation for the universality in glasses

There are still many fundamental questions regarding glasses  
to think about and work on



Happy Birthday