## Wave Equation

The wave equation

$$
\begin{equation*}
\frac{\partial^{2} \phi}{\partial x^{2}}-\frac{1}{c^{2}} \frac{\partial^{2} \phi}{\partial t^{2}}=0 \tag{1}
\end{equation*}
$$

with the initial values (Cauchy data) $\phi(x, 0), \dot{\phi}(x, 0)$ is most easily solved through factorization

$$
\begin{equation*}
\left(\frac{\partial}{\partial x}+\frac{1}{c} \frac{\partial}{\partial t}\right)\left(\frac{\partial}{\partial x}-\frac{1}{c} \frac{\partial}{\partial t}\right) \phi=0 \tag{2}
\end{equation*}
$$

However, historically the first people to solve (1) had slightly different interpretations of (2). Euler interpreted it geometrically while d'Alembert interpreted it algebraically.

## Euler's solution

Euler wants to see what the left most operator in (2) does geometrically. He realized that if we can find a parametrization $x\left(s_{1}\right), t\left(s_{1}\right)$ such that

$$
\left\{\begin{array} { l } 
{ \frac { \partial x } { \partial s _ { 1 } } = 1 }  \tag{3}\\
{ \frac { \partial t } { \partial s _ { 1 } } = \frac { 1 } { c } }
\end{array} \Rightarrow \left\{\begin{array}{l}
x=s_{1}+a \\
t=\frac{1}{c} s_{1}+b
\end{array}\right.\right.
$$

then we may rewriting (2) as

$$
\begin{equation*}
\frac{\partial}{\partial s_{1}} A\left(s_{1}\right)=0 \tag{4}
\end{equation*}
$$

Voila, the operator $\frac{\partial}{\partial x}+\frac{1}{c} \frac{\partial}{\partial t}$ does not change the value of its operand along the curve $\gamma\left(s_{1}\right)=\left(x\left(s_{1}\right), t\left(s_{1}\right)\right)$. Therefore we can calculate the value of its operand at any point ( $x, y$ ) by tracing back to the Cauchy data, indeed, using $s_{1}^{\prime}=-c b$ we have

$$
\begin{array}{r}
A\left(x\left(s_{1}^{\prime}\right), t\left(s_{1}^{\prime}\right)\right)=A(a-c b, 0)=A\left(x\left(s_{1}\right)-c t\left(s_{1}\right), 0\right) \Rightarrow \\
A(x, t)=A(x-c t, 0) \tag{5}
\end{array}
$$

where I have obtained from (3) that

$$
\left\{\begin{array}{l}
a=x-s_{1}  \tag{6}\\
b=t-\frac{1}{c} s_{1}
\end{array}\right.
$$

Reverse the order of factorization of (2), we can define $s_{2}$ such that

$$
\begin{equation*}
\frac{\partial}{\partial s_{2}} B\left(s_{2}\right)=0 \tag{7}
\end{equation*}
$$

with

$$
\begin{equation*}
B(x, t)=B(x+c t, 0) \tag{8}
\end{equation*}
$$

Writing out $A, B$ explicitly

$$
\left\{\begin{array}{l}
A=\phi^{\prime}-\frac{1}{c} \dot{\phi}  \tag{9}\\
B=\phi^{\prime}+\frac{1}{c} \dot{\phi}
\end{array}\right.
$$

we have

$$
\begin{align*}
\dot{\phi}(x, t) & =\frac{c}{2}\{B(x, t)-A(x, t)\} \\
& =\frac{c}{2}\{B(x+c t, 0)-A(x-c t, 0)\} \tag{10}
\end{align*}
$$

which is a known function. The solution can therefore obtained by integrating over time

$$
\begin{equation*}
\phi(x, T)=\phi(x, 0)+\int_{0}^{T} \dot{\phi}(x, t) d t \tag{11}
\end{equation*}
$$

This is a non-trivial integral, but can be done with lots of change of variable (which results in change of integration limits) and the final form is

$$
\begin{equation*}
\phi(x, T)=\frac{1}{2}\{\phi(x+c t, 0)+\phi(x-c t, 0)\}+\frac{1}{2 c} \int_{x-c t}^{x+c t} \dot{\phi}(\xi, 0) d \xi \tag{12}
\end{equation*}
$$

Euler's solution gives nice geometrical interpretations to the operators and shows that the final solution only depends on Cauchy data that can be traced back within a light cone. However, it is mathematically tedious.

## d'Alembert's solution

d'Alembert starts by changing to the light-cone coordinates

$$
\left\{\begin{array} { l } 
{ \xi = x + c t }  \tag{13}\\
{ \eta = x - c t }
\end{array} \Rightarrow \left\{\begin{array}{l}
x=\frac{1}{2}(\xi+\eta) \\
c t=\frac{1}{2}(\xi-\eta)
\end{array}\right.\right.
$$

this way we may rewrite (2) as

$$
\begin{equation*}
4 \frac{\partial^{2}}{\partial \xi \partial \eta} \phi(\xi, \eta)=0 \tag{14}
\end{equation*}
$$

the general solution is therefore

$$
\begin{equation*}
\phi(\xi, \eta)=f(\xi)+g(\eta) \tag{15}
\end{equation*}
$$

or

$$
\begin{equation*}
\phi(x, t)=f(x+c t)+g(x-c t) \tag{16}
\end{equation*}
$$

now the initial values say

$$
\begin{gather*}
\left\{\begin{array}{l}
\phi(x, 0)=f(x)+g(x) \\
\dot{\phi}(x, 0)=c\left(f^{\prime}(x)-g^{\prime}(x)\right)
\end{array} \Rightarrow\right. \\
\left\{\begin{array}{l}
f(x)+g(x)=\phi(x, 0) \\
f(x)-g(x)=\frac{1}{c} \int_{0}^{x} \dot{\phi}(\xi, 0) d \xi+C
\end{array} \Rightarrow\right. \\
\left\{\begin{array}{l}
f(x)=\frac{1}{2} \phi(x, 0)+\frac{1}{2 c} \int_{0}^{x} \dot{\phi}(\xi, 0) d \xi+\frac{1}{2} C \\
g(x)=\frac{1}{2} \phi(x, 0)-\frac{1}{2 c} \int_{0}^{x} \dot{\phi}(\xi, 0) d \xi-\frac{1}{2} C
\end{array}\right. \tag{17}
\end{gather*}
$$

and we have obtained our solution

$$
\begin{equation*}
\phi(x, t)=\frac{1}{2}\{\phi(x+c t, 0)+\phi(x-c t)\}+\frac{1}{2 c} \int_{x-c t}^{x+c t} \dot{\phi}(\xi, 0) d \xi \tag{18}
\end{equation*}
$$

Notice, this is the same as Euler's solution, but the steps are a lot more clearly laid out. It seems a bit like cheating though since the light-cone coordinates came out of nowhere. (I'm pretty sure he got the inspiration from Euler)

