

# Functional Derivative

The goal of this section is to discover a suitable definition of a "functional derivative", such that we can take the derivative of a functional and still have the same rules of differentiation as normal calculus. For example, we wish to find a definition for  $\frac{\delta J}{\delta y}$ , where  $J[y(x)]$  is a functional of  $y(x)$  such that things like  $\frac{\delta}{\delta y} J^2 = 2J \frac{\delta J}{\delta y}$  are still true.

## Definitions

### Functional

Stone's definition of *local functional* where  $f$  is a multivariable function

$$J[y] = \int_{x_1}^{x_2} f(x, y(x), y'(x), y''(x), \dots, y^{(n)}(x)) dx = \int_{x_1}^{x_2} f dx \quad (1)$$

Notice the functional  $J$  "eats" an entire function  $y$ , which is defined using its local values  $y(x), y'(x)$  etc, and spits out a number through integration. In short, a functional is just a number that depends on an input function.

### Variation

A *variation* of the functional is the amount the functional changes when the input function is changed by a little bit. Suppose we let  $y(x) \rightarrow y(x) + \delta y(x)$ , then since  $\frac{d}{dx}$  is linear

$$\begin{cases} y'(x) \rightarrow y'(x) + \frac{d}{dx} \delta y(x) = y'(x) + \delta y'(x) \\ y''(x) \rightarrow y''(x) + \frac{d^2}{dx^2} \delta y(x) = y''(x) + \delta y''(x) \\ \vdots \\ y^{(n)}(x) \rightarrow y^{(n)}(x) + \frac{d^n}{dx^n} \delta y(x) = y^{(n)}(x) + \delta y^{(n)}(x) \end{cases} \quad (2)$$

thus the integrand of the new output  $J[y + \delta y]$  can be expanded to first order using Taylor expansion of a multivariable function around the old input  $y$

$$\begin{aligned} J[y + \delta y] &= \int_{x_1}^{x_2} f(x, y + \delta y, y' + \delta y', \dots, y^{(n)} + \delta^{(n)} y) dx \\ &= \int_{x_1}^{x_2} \left\{ f + \frac{\partial f}{\partial x} dx + \frac{\partial f}{\partial y} \delta y + \frac{\partial f}{\partial y'} \delta y' + \dots + \frac{\partial^{(n)} f}{\partial y^{(n)}} \delta y^{(n)} \right\} dx \quad (3) \end{aligned}$$

The variation of the functional is thus, by definition, the new output minus the old output taken to first order.

$$\begin{aligned}\delta J &= J[y + \delta y] - J[y] \\ &= \int_{x_1}^{x_2} \left\{ \frac{\partial f}{\partial x} dx + \frac{\partial f}{\partial y} \delta y + \frac{\partial f}{\partial y'} \delta y' + \cdots + \frac{\partial^{(n)} f}{\partial y^{(n)}} \delta y^{(n)} \right\} dx\end{aligned}\quad (4)$$

we can moved all the  $\frac{d}{dx}$  on  $\delta y$  to  $f$  using integration by parts

$$\begin{aligned}\delta J &= \frac{d}{dx} \left( \frac{\partial f}{\partial y'} \right) \delta y(x) \Big|_{x_1}^{x_2} + \frac{d}{dx} \left( \frac{\partial f}{\partial y''} \right) \delta y'(x) \Big|_{x_1}^{x_2} - \frac{d^2}{dx^2} \left( \frac{\partial f}{\partial y''} \right) \delta y(x) \Big|_{x_1}^{x_2} \\ &+ \cdots + (-1)^{n-1} \frac{d^n}{dx^n} \left( \frac{\partial f}{\partial y^{(n)}} \right) \delta y^{(n)}(x) \Big|_{x_1}^{x_2} + \int_{x_1}^{x_2} \frac{\partial f}{\partial x} dx + \\ &\int_{x_1}^{x_2} \left( \frac{\partial f}{\partial y} - \frac{d}{dx} \frac{\partial f}{\partial y'} + \frac{d^2}{dx^2} \frac{\partial f}{\partial y''} - \cdots + (-1)^{n-1} \frac{d^n}{dx^n} \frac{\partial f}{\partial y^{(n)}} \right) \delta y(x) dx\end{aligned}\quad (5)$$

and Voila! (5) is the *variation* of the local functional defined in (1) with no additional assumption.

## Functional Derivative

For a normal multi-variable function  $f(x_1, x_2, \dots, x_n)$  we have a nice form for its variation

$$df = \sum_{i=1}^n \left\{ \frac{\partial f}{\partial x_i} dx_i \right\}\quad (6)$$

and we know how to calculate the derivatives  $\frac{\partial f}{\partial x_i}$ . Here we wish to rewrite (5) such that we have a similar form for the variation of a functional

$$\delta J = \int_{x_1}^{x_2} dx \left\{ \frac{\delta J}{\delta y}(x) \delta y(x) \right\}\quad (7)$$

Unfortunately, this is only possible under special circumstances. That is, we need the variation to have "**fixed-ends**" ( $\delta y^{(n)}(x_1) = \delta y^{(n)}(x_2) = 0$ ) and that we require **implicit f** ( $\frac{\partial f}{\partial x} = 0$ )<sup>1</sup>. Basically, we want everything before the last line of (5) to vanish. This way, comparing (5) and (7) we finally have

$$\boxed{\frac{\delta J}{\delta y} = \frac{\partial f}{\partial y} - \frac{d}{dx} \frac{\partial f}{\partial y'} + \frac{d^2}{dx^2} \frac{\partial f}{\partial y''} - \cdots + (-1)^{n-1} \frac{d^n}{dx^n} \frac{\partial f}{\partial y^{(n)}}}\quad (8)$$

<sup>1</sup>This is slightly over kill, since we just want  $\int_{x_1}^{x_2} \frac{\partial f}{\partial x} dx = 0$

As long as  $\delta J$  can be written as (7), we will have our nice rules of differentiation. For example

$$\begin{aligned}\delta(J^2) &\equiv (J + \delta J)^2 - J^2 = 2J\delta J + O(\delta J^2) \\ &= \int_{x_1}^{x_2} dx \left\{ 2J \frac{\delta J}{\delta y}(x) \delta y(x) \right\} \Rightarrow \\ \frac{\delta J^2}{\delta y} &= 2J \frac{\delta J}{\delta y}\end{aligned}\tag{9}$$

A word of caution: This definition of functional derivative is nice, but as  $f$  involves higher derivatives of  $y$ , the fixed-end condition becomes harsher and the range of  $y$  this derivative applies to quickly diminishes. Therefore it is sometimes more useful to make variations by hand according to (5).

## Lagrangian Mechanics

When the integrand of the functional only has dependence on  $y$  and  $y'$  ( $f(y, y')$ ), (8) reduces to the popular *Fréchet derivative*

$$\frac{\delta J}{\delta y} = \frac{\partial f}{\partial y} - \frac{d}{dx} \frac{\partial f}{\partial y'}\tag{10}$$

this form should look familiar to all physicists, since it laid the foundation for basic Lagrangian mechanics. In a typical classical mechanics problem, we wish to minimize the *action*  $S$ , which is often a functional of a *configuration function*  $q$ , whose basic *independent variable* is time  $t$ . That is

$$S[q] = \int_{t_1}^{t_2} L(q(t), \dot{q}(t)) dt\tag{11}$$

where  $L$  is the *Lagrangian*. To find extrema, we set derivative to 0

$$\begin{aligned}\frac{\delta S}{\delta q} &= 0 \Rightarrow \\ \frac{\partial L}{\partial q} &= \frac{d}{dt} \frac{\partial L}{\partial \dot{q}}\end{aligned}\tag{12}$$

Lo and behold, the *Lagrangian equation of motion*.