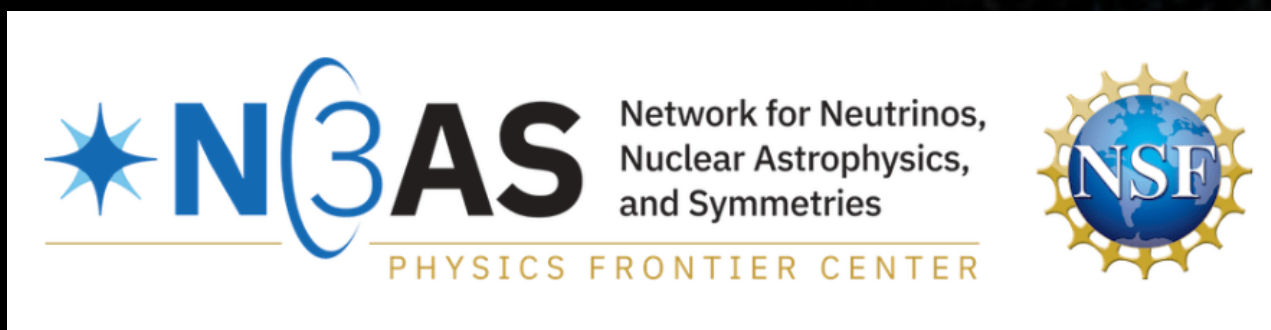


DARK MATTER AND GRAVITATIONAL WAVE SCIENCE

BEYOND THE STANDARD MODEL WITH GRAVITATIONAL WAVE OBSERVATIONS

DIVYA SINGH
UNIVERSITY OF CALIFORNIA, BERKELEY





Bertone et al. (2019)

This is a rich and exciting field of research with great avenues for multi-messenger detection of dark matter candidates!



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**Compact binary
Coalescences and
Dark matter**

DETECTING GRAVITATIONAL WAVES FROM CBC(S)

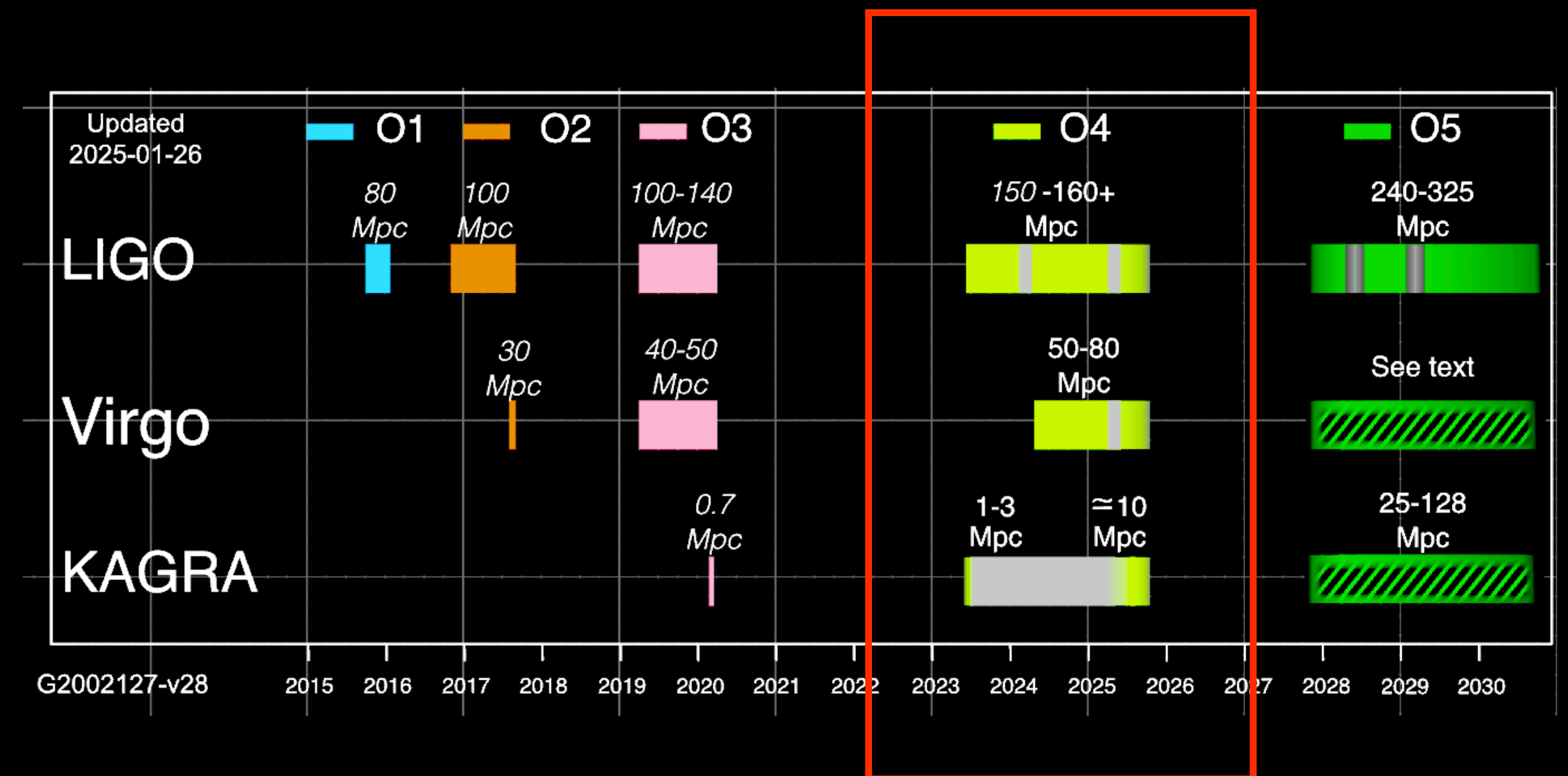
- Real-time searches look for BBH, BNS, and NSBH signals with multiple search pipelines.
- >200* public alerts from significant detections to date in the fourth observing run of LIGO-Virgo-Kagra.
- Total detections so far: >290

First BBH (O1): GW150914 PBHs? Sasaki et al. (2016)

First BNS (O2): GW170817

First NSBHs (O3): GW200105, GW200115

NSBH in O4: GW230529 (Mass Gap BH)

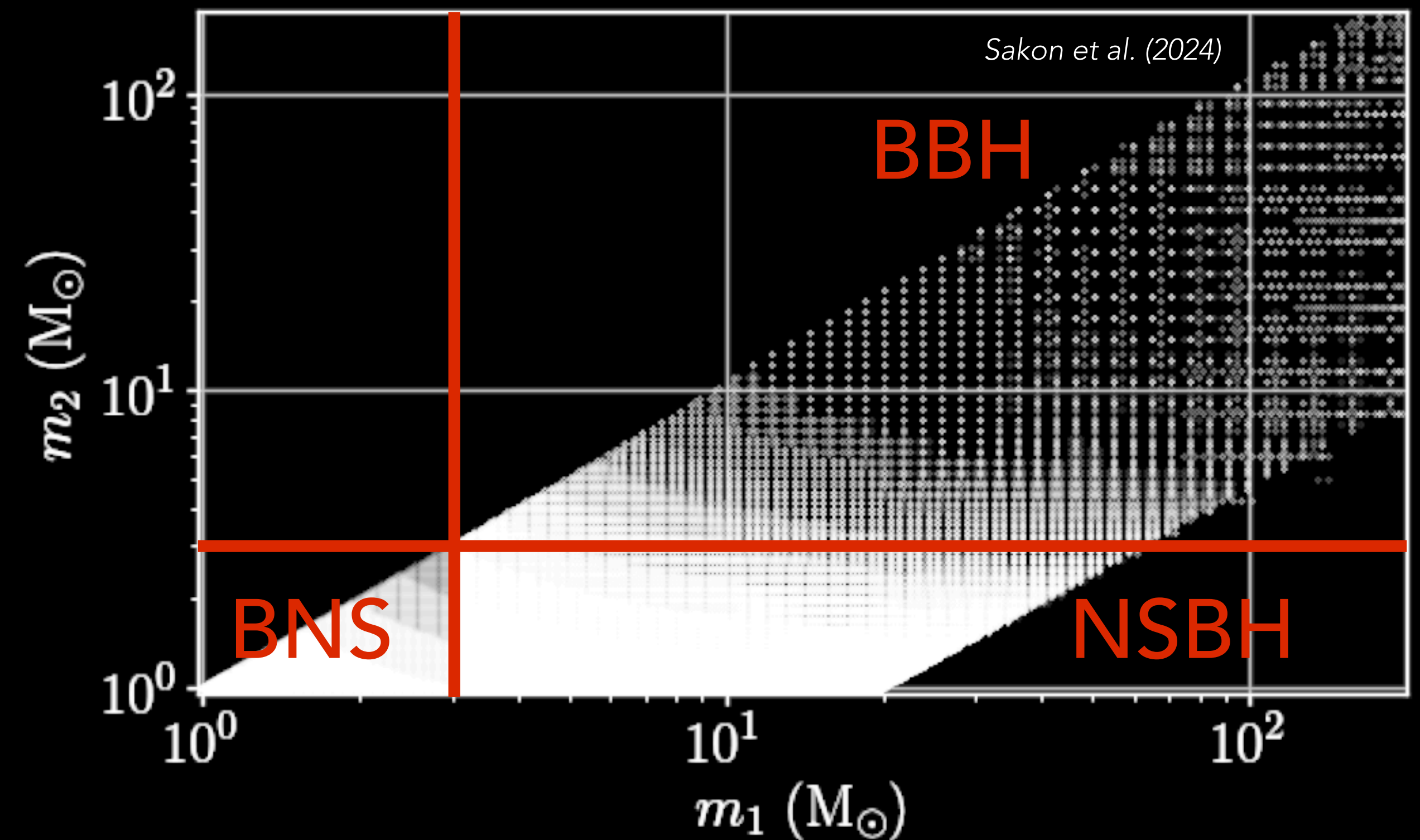


From: <https://emfollow.docs.ligo.org/userguide/capabilities.html>

DETECTING GRAVITATIONAL WAVES FROM CBC(S)

- Real-time searches look for BBH, BNS, and NSBH signals with multiple search pipelines.
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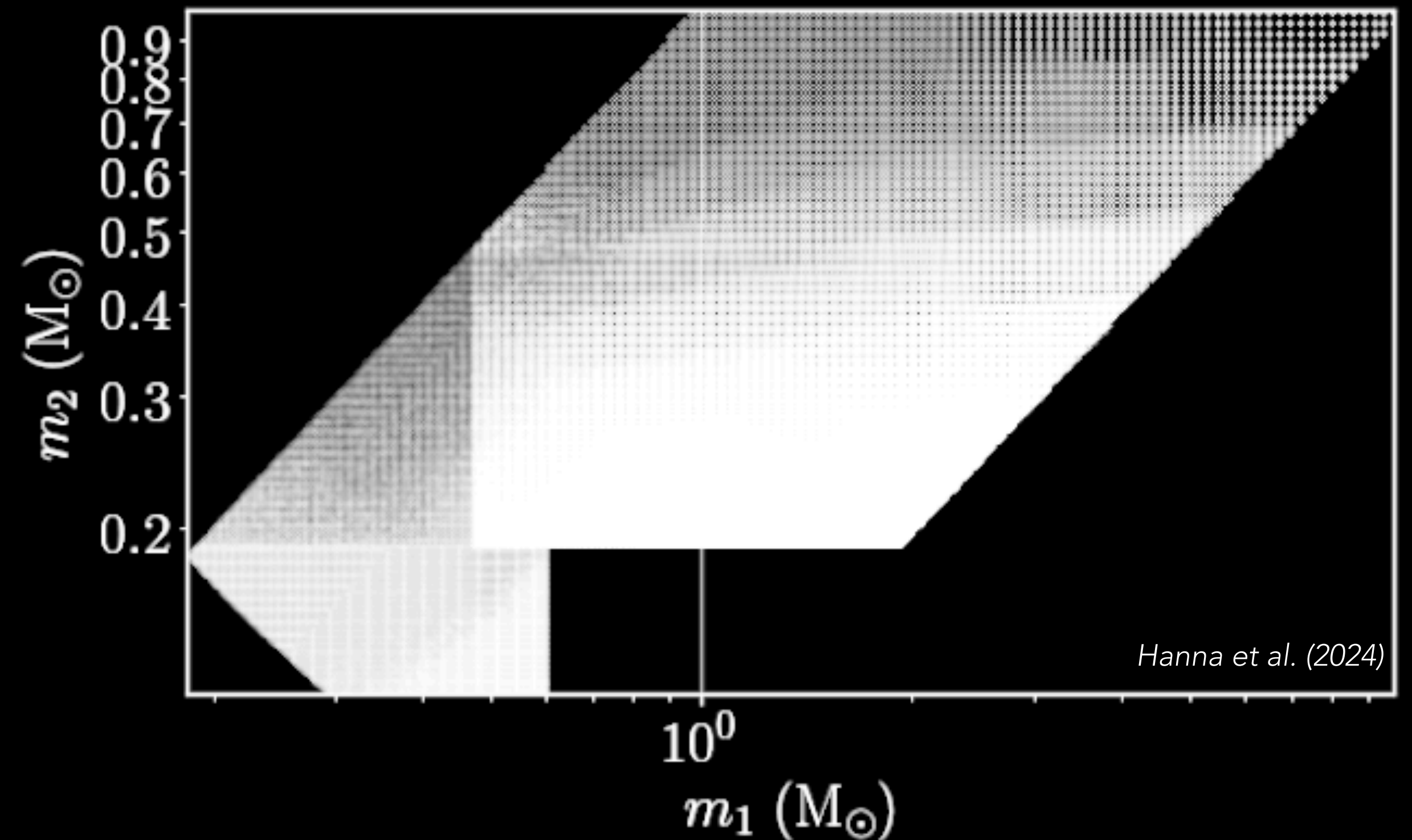
Mass of Compact objects from stellar evolution $> 1.4 M_{\odot}$



FORMATION CHANNELS BEYOND STELLAR EVOLUTION

SUB-SOLAR MASS SEARCHES IN GRAVITATIONAL WAVE DATA

- Primordial Black Holes
- Dissipative dark matter — Dark Black Holes
- Neutron Star implosion from dark matter accumulation

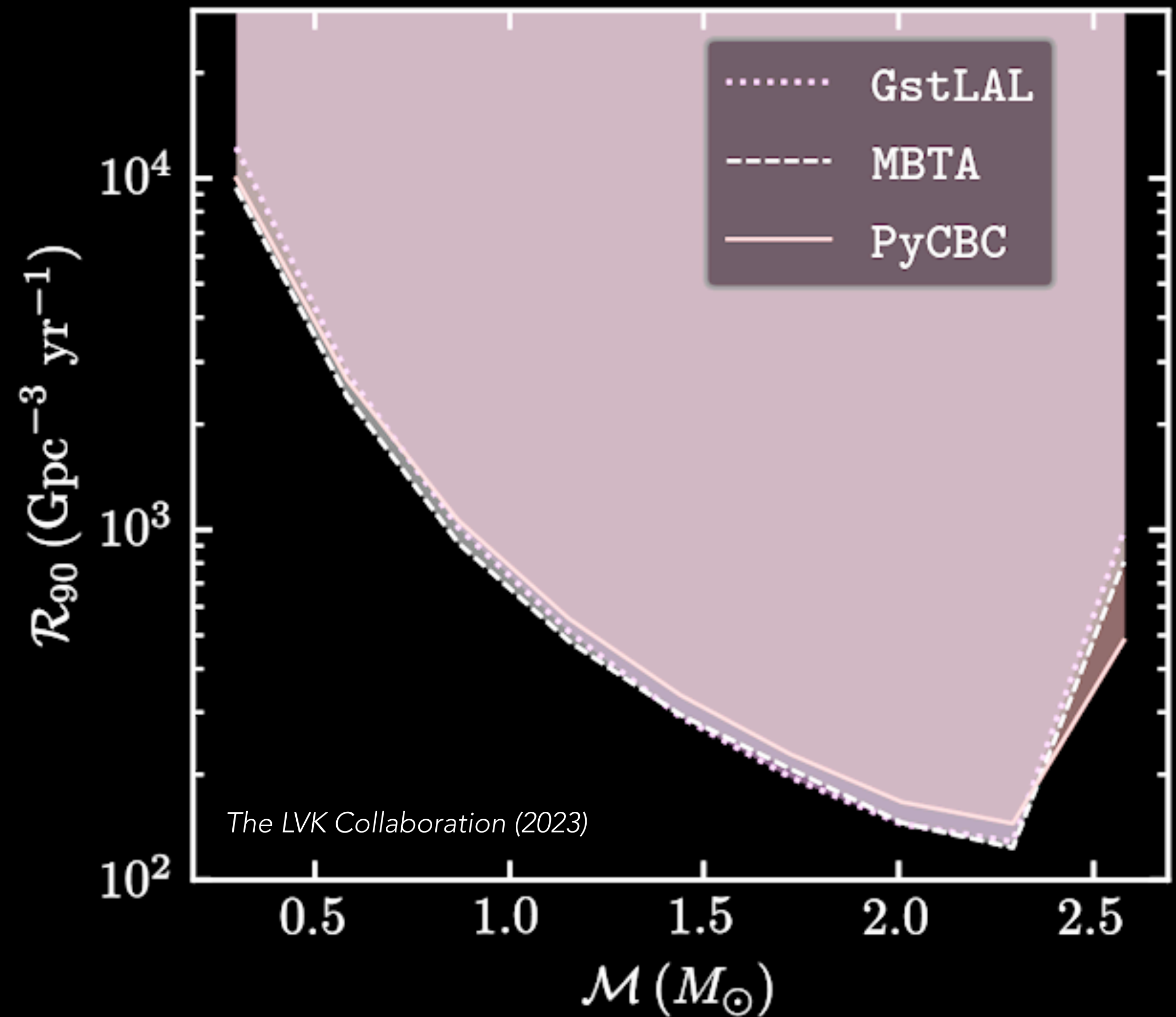


SUB-SOLAR MASS SEARCHES IN DATA FROM LIGO-VIRGO

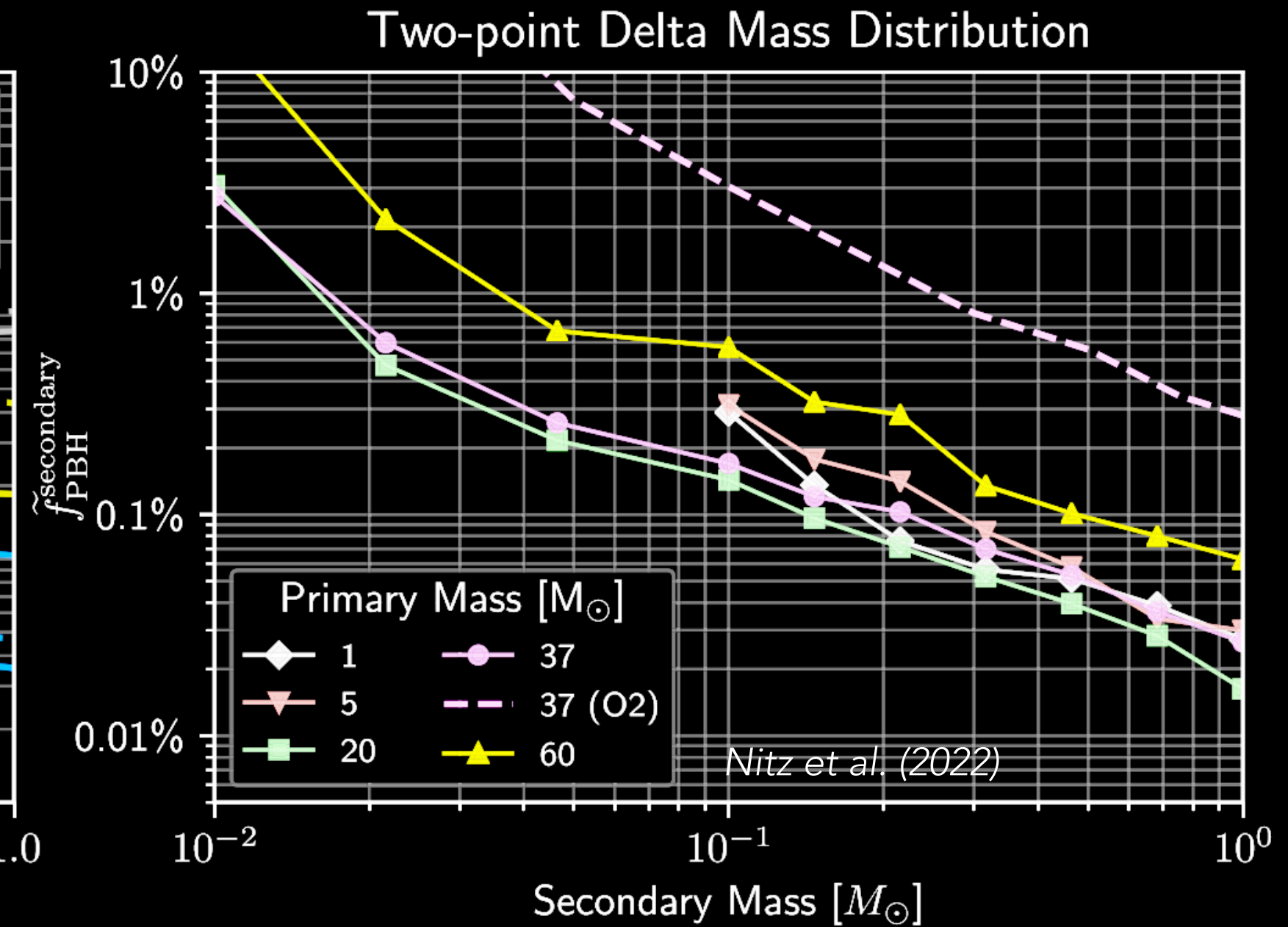
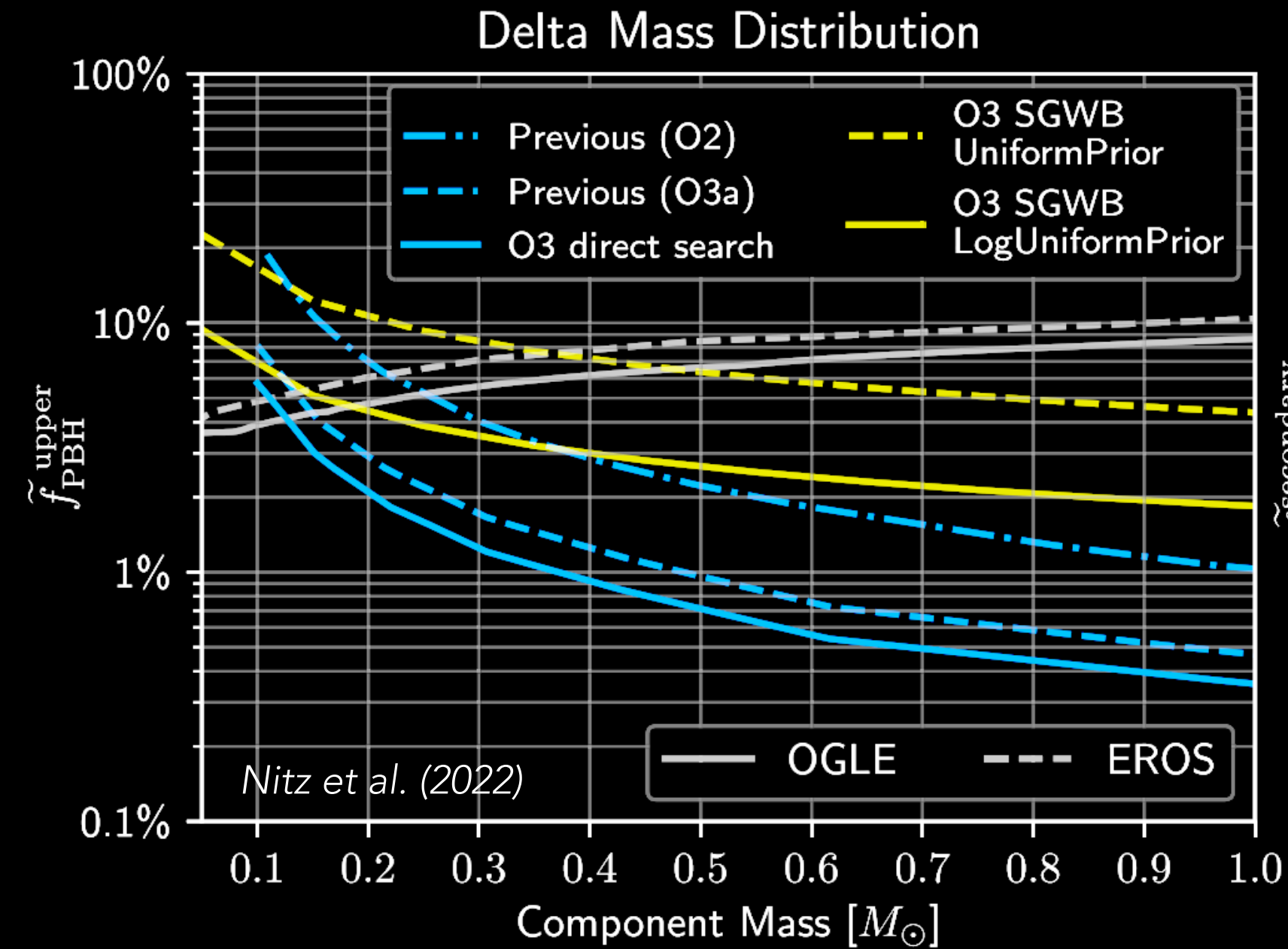
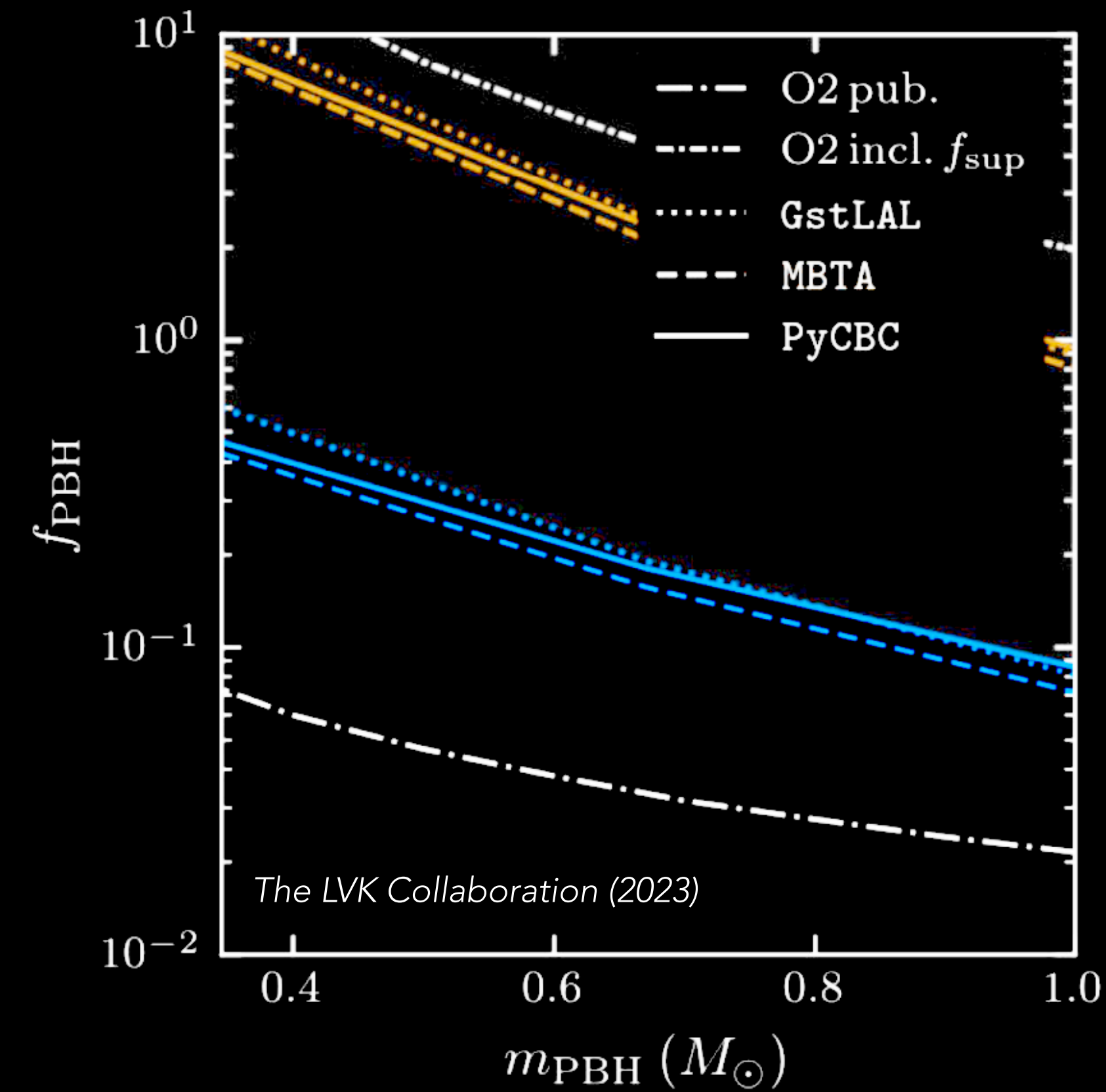
- Searches performed by multiple independent search pipelines.

Nitz et al. (2022), The LVK Collaboration (2023)

- No detections reported from searches in data from O1, O2 and O3.
- Upper limits on Merger Rates
- $100 \text{ per Gpc}^3 \text{ yr}^{-1}$ at chirp mass of $2 M_{\text{sun}}$.



PRIMORDIAL BLACK HOLES



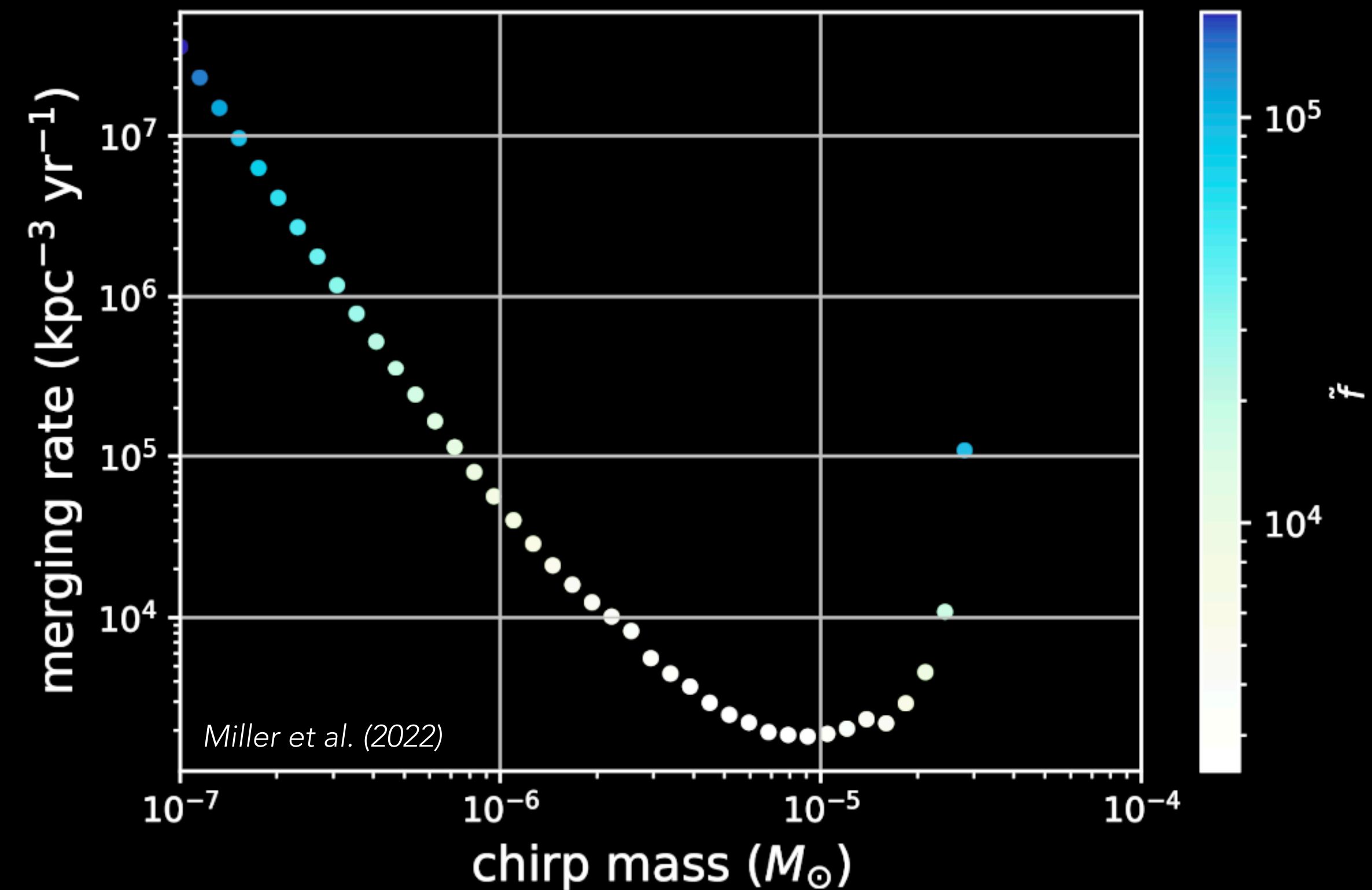
Constraints on the abundance of sub-solar mass PBHs
From astrophysical rate limits.

PRIMORDIAL BLACK HOLES

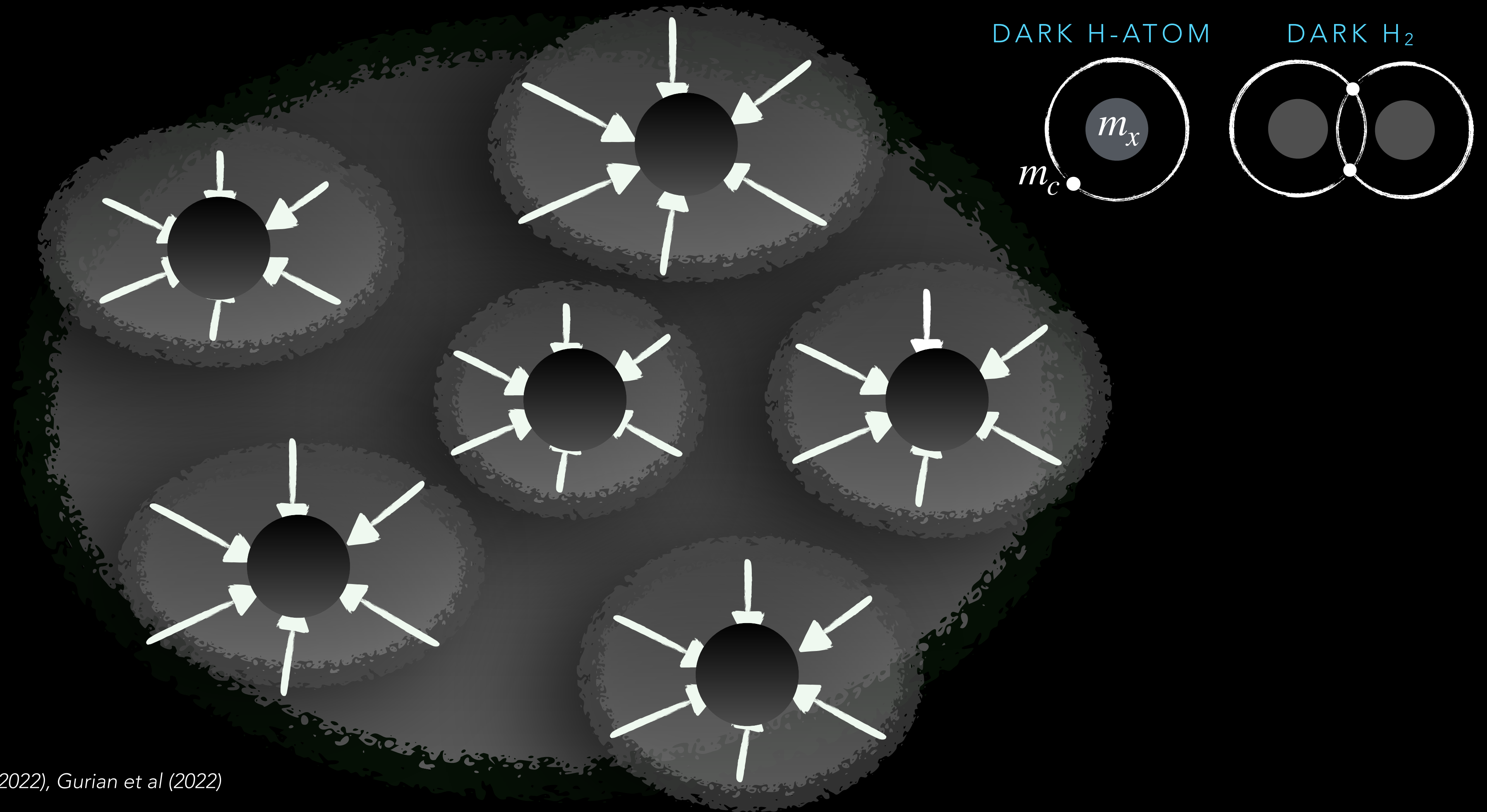
BEYOND SUB-SOLAR MASS SEARCHES IN GW STRAIN DATA

Quasi-monochromatic GW signals from **planetary-mass and asteroid-mass** PBH binaries.

Use limits from continuous wave searches to constrain the merger rate and the abundance of PBH binaries.



DISSIPATIVE DARK MATTER

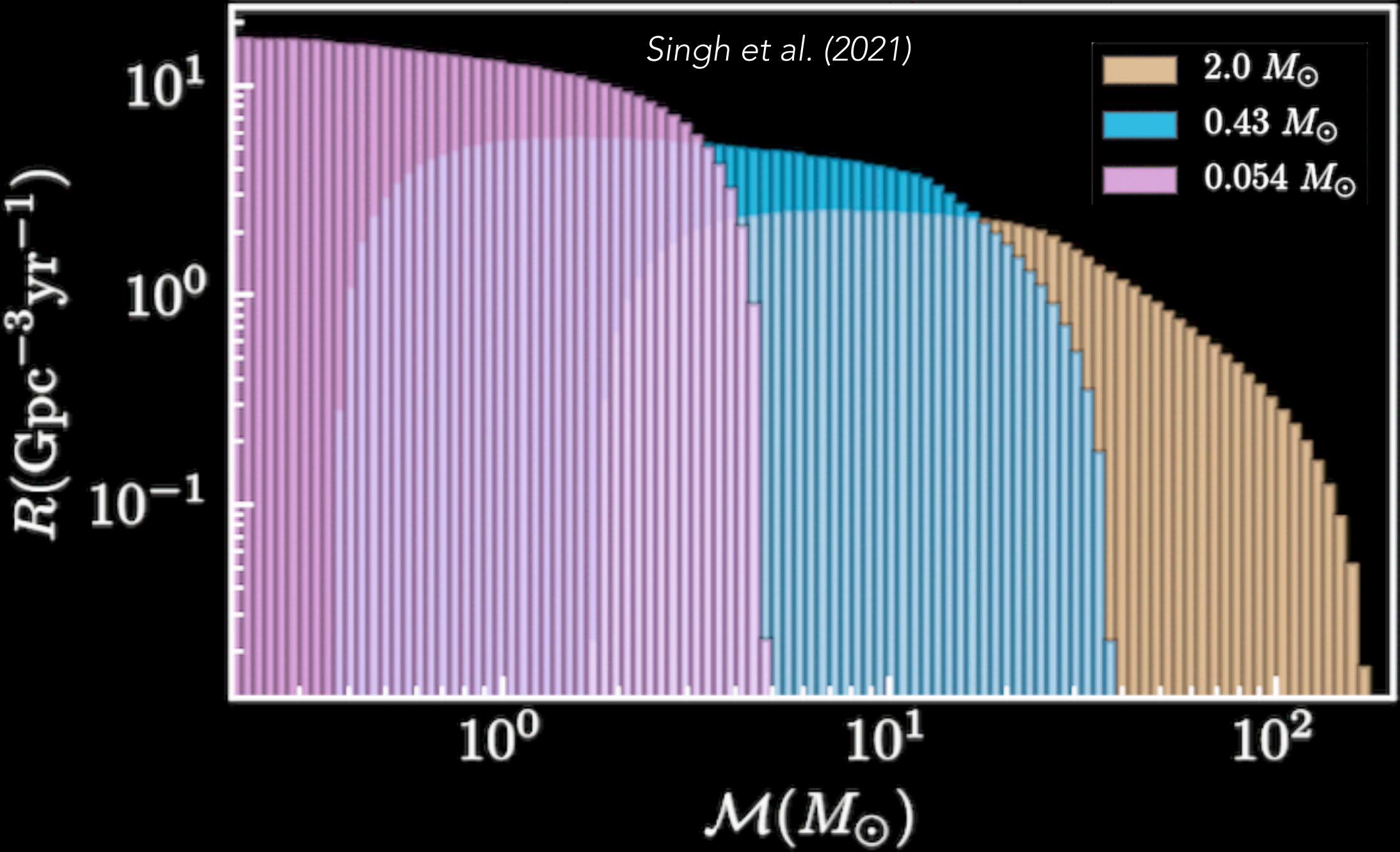


Shandera et al. (2018), Ryan et al (2022), Gurian et al (2022)

DISSIPATIVE DARK MATTER

Possibly the minimum mass of dark BHs

$$M_{\text{Chandrasekhar}}^{\text{Dark}} = 1.4 M_{\odot} \left(\frac{m_p}{m_x} \right)^2$$



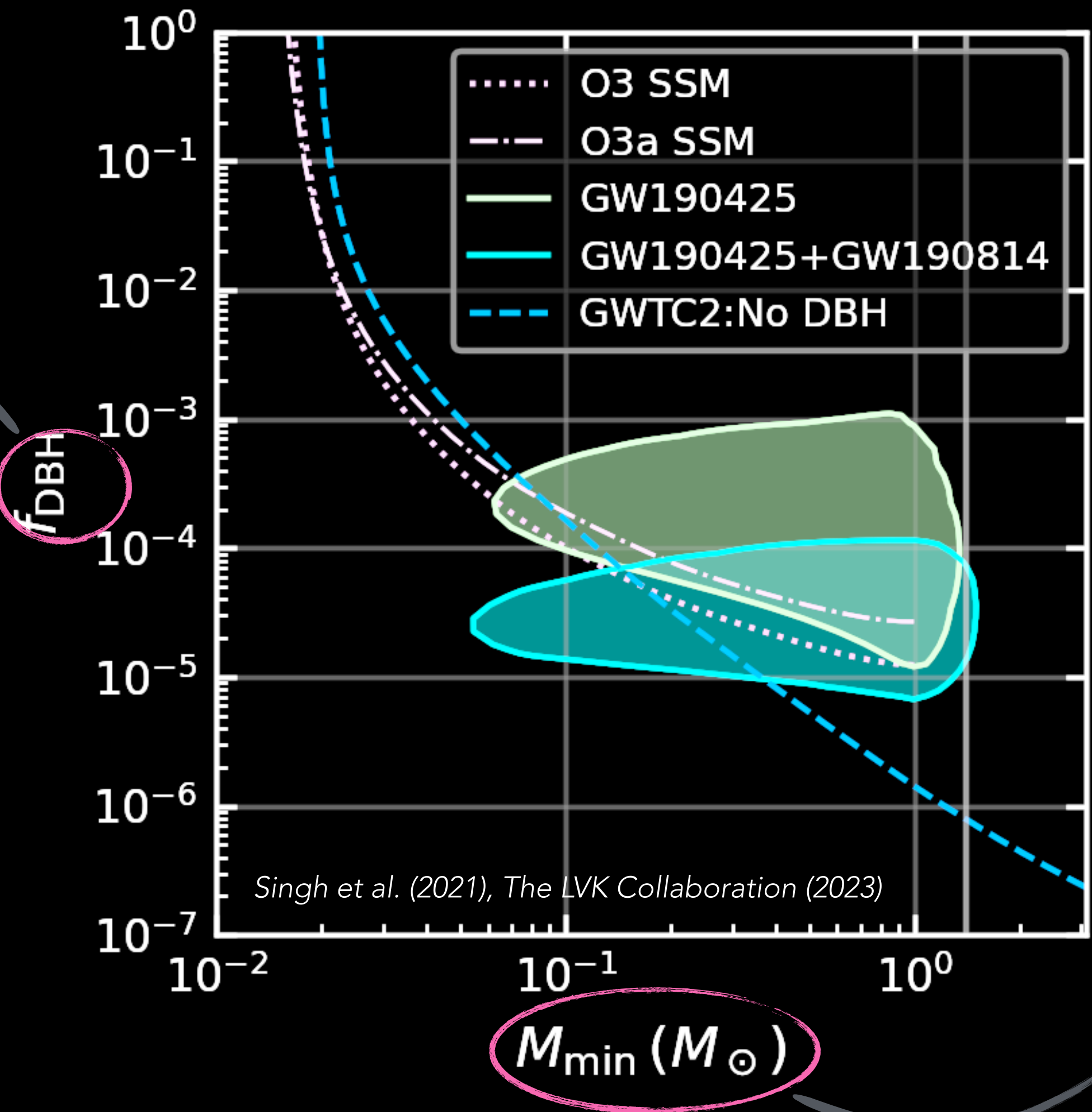
m_X [GeV]	m_c [keV]	$M_{\text{Chand.}}^{\text{dark}}$ [$10^{-5} M_\odot$]	M_{DBH} [M_\odot]	Rates per year				$m_1 < 1.4$ [%]	$m_1, m_2 < 1.4$ [%]
				raw (MWEG $^{-1}$)	aLIGO (current)	aLIGO (full)	Einstein T.		
62	31	33	0.0068 – 0.68	$2.0 \times 10^{-6} (10^{-4})$	0.0012 (0.12)	0.020 (2.0)	60 (6000)	100%	100%
48	47	56	0.016 – 1.6	$1.3 \times 10^{-6} (10^{-4})$	0.0065 (0.65)	0.11 (11)	330 (33k)	99%	79%
32	70	125	0.054 – 5.4	$6.6 \times 10^{-7} (10^{-5})$	0.068 (6.8)	1.1 (110)	3500 (350k)	53%	9.3%
16	140	500	0.43 – 43	$1.9 \times 10^{-7} (10^{-5})$	0.89 (89)	22 (2200)	92k (9200k)	9.8%	0.14%

TABLE I. DBH masses and binary merger rates today, estimated using the procedure in the text, for several choices of dark proton mass m_X and dark electron mass m_c . All black hole masses are given in solar masses. In all cases we have set the dark fine structure constant to $\alpha_D = 0.01$ and the ratio of present day temperature of the dark sector to photon temperature to $\xi = 0.02$. The conservative (optimistic) rates use $f_{\text{cool}} \times f_{\text{form. eff.}} = 10^{-5} (10^{-3})$. The optimistic rate for $m_X = 16$ GeV is high enough it may already be constrained by existing LIGO data. The last two columns show the percent of binaries where one or both black holes in the binary has a mass less than the standard Chandrasekhar mass ($1.4 M_\odot$).

Appreciable merger rates albeit
dependent on dark matter particle masses

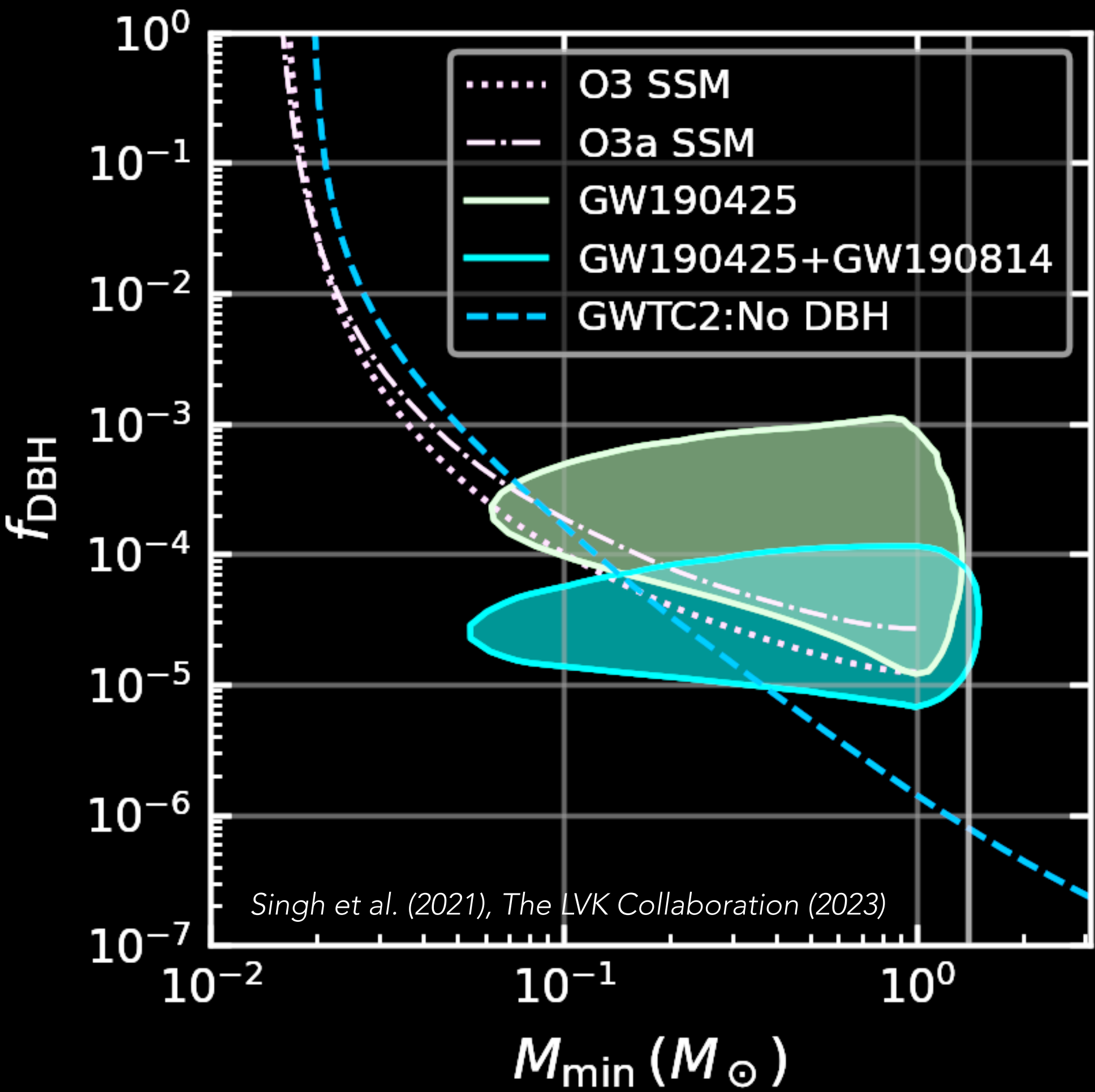
DISSIPATIVE DARK MATTER

Abundance of dark black holes - Fraction of DM in dark black holes



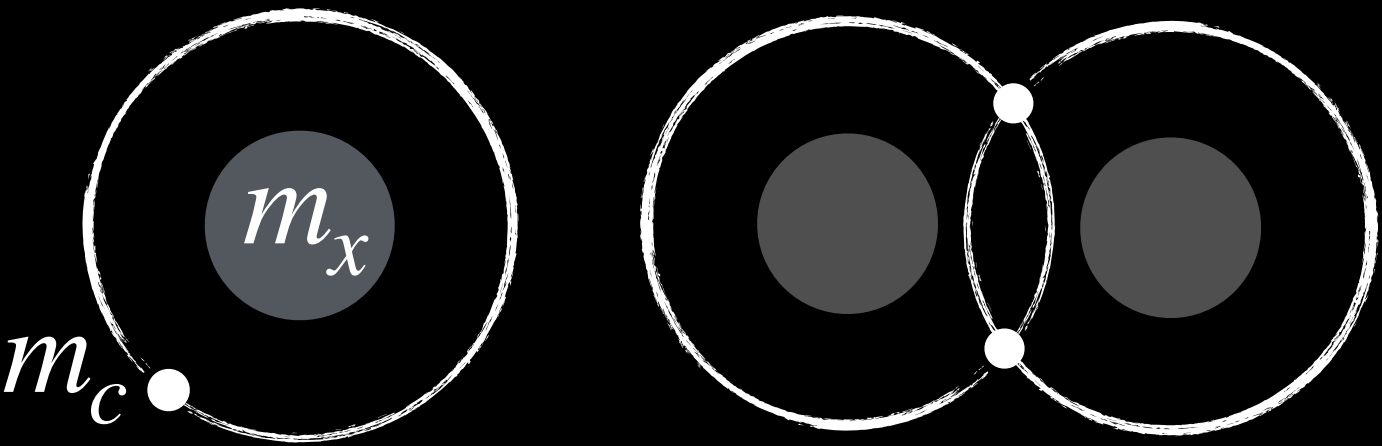
Minimum possible mass dark black holes

DISSIPATIVE DARK MATTER



DARK H-ATOM

DARK H₂



Possibly the minimum mass of dark BHs

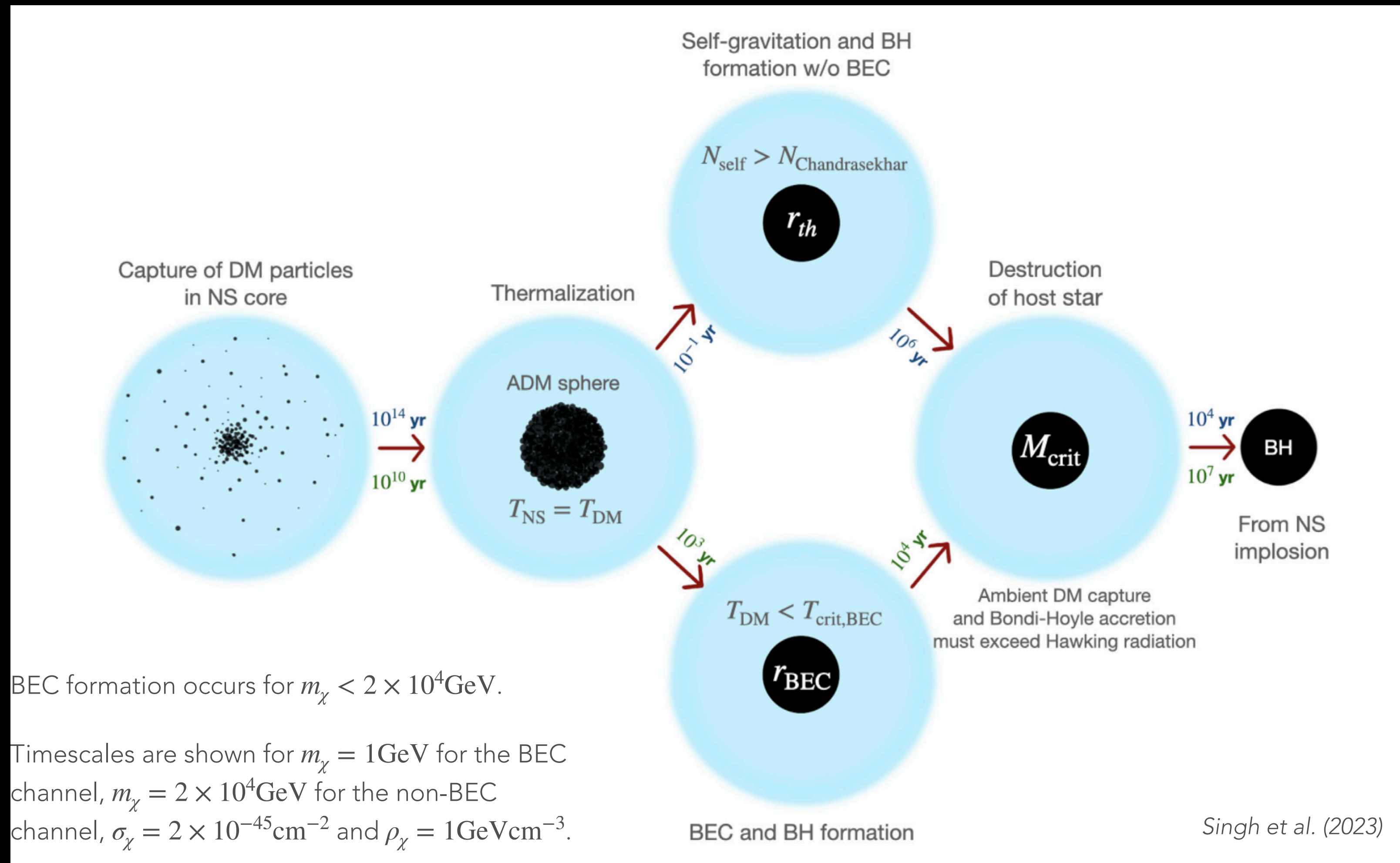
$$M_{\text{Chandrasekhar}}^{\text{Dark}} = 1.4M_{\odot} \left(\frac{m_p}{m_x} \right)^2$$

TABLE I. Probable minimum masses of dark black holes, M_{min} , and the corresponding heavy fermion masses, m_x , for the two cases of observed dark black-hole binaries. Heavy fermion masses are computed using the dark matter Chandrasekhar mass limit, by setting $M_{\text{DC}} = M_{\text{min}}$ determined from the data.

Observed DBH binary	\mathcal{M}/M_{\odot}	M_{min}/M_{\odot}	m_x/GeV
GW190425	1.44	0.062–1.34	0.95–4.44
GW190425, GW190814	1.44, 6.1	0.054–1.50	0.91–4.76

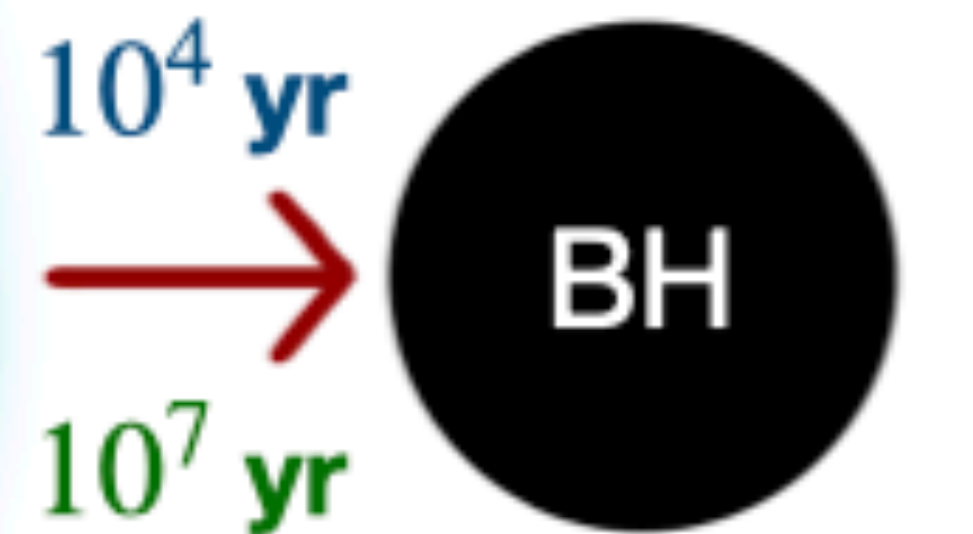
ASYMMETRIC DARK MATTER

BLACK HOLES FROM NEUTRON STAR IMPLOSION



Black holes with neutron star like masses: $1-3 M_\odot$

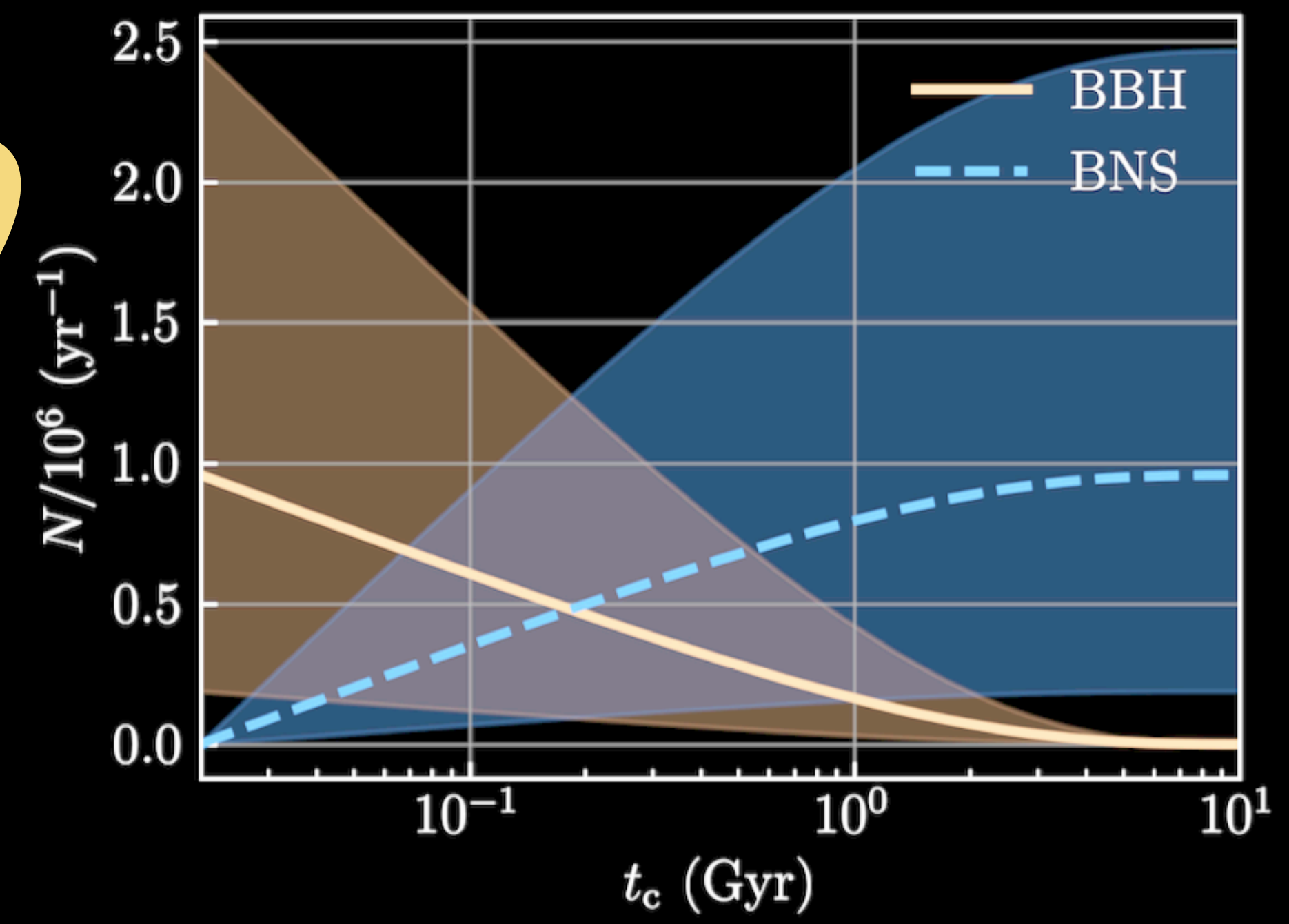
on
ar



From NS
implosion

capture
e accretion
king radiation

Time to collapse > Merger time



Time to collapse < Merger time

ASYMMETRIC DARK MATTER

BLACK HOLES FROM NEUTRON STAR IMPLOSION

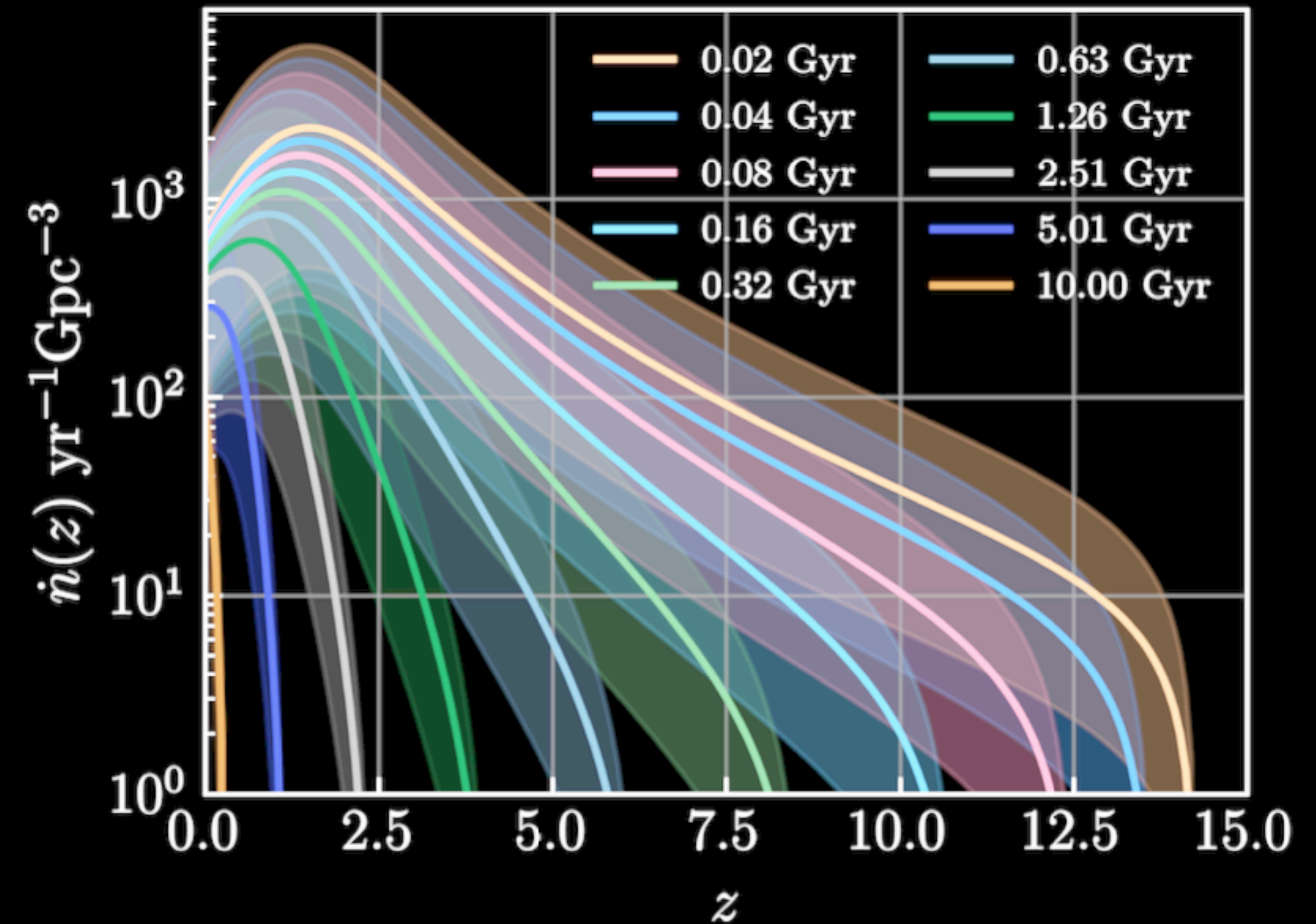
$$\dot{n}(z) = A \int_{t_d^{\min}}^{t_d^{\max}} \underbrace{\psi(z_f(z, t_d))}_{\text{Star formation rate}} \underbrace{\mathcal{P}(t_d)}_{\text{Delay time distribution}} dt_d$$

$$\dot{n}(z) = \dot{n}(z)_{\text{BNS}} + \dot{n}(z)_{\text{BBH}}$$

collapse time $\int_{t_d^{\min}}^{\text{collapse time}}$

collapse time $\int_{\text{collapse time}}^{t_d^{\max}}$

Merger rate density for BBH is a function of the collapse time. Higher collapse time implies more BNS and vice versa.



ASYMMETRIC DARK MATTER

BLACK HOLES FROM NEUTRON STAR IMPLOSION

Detectable
~~TOTAL MERGER~~
RATE

$$\dot{N} = \int_0^z \frac{\dot{n}(z')}{1+z'} \frac{dV_c}{dz'} dz' \epsilon(z') \rightarrow \text{Detector Efficiency}$$

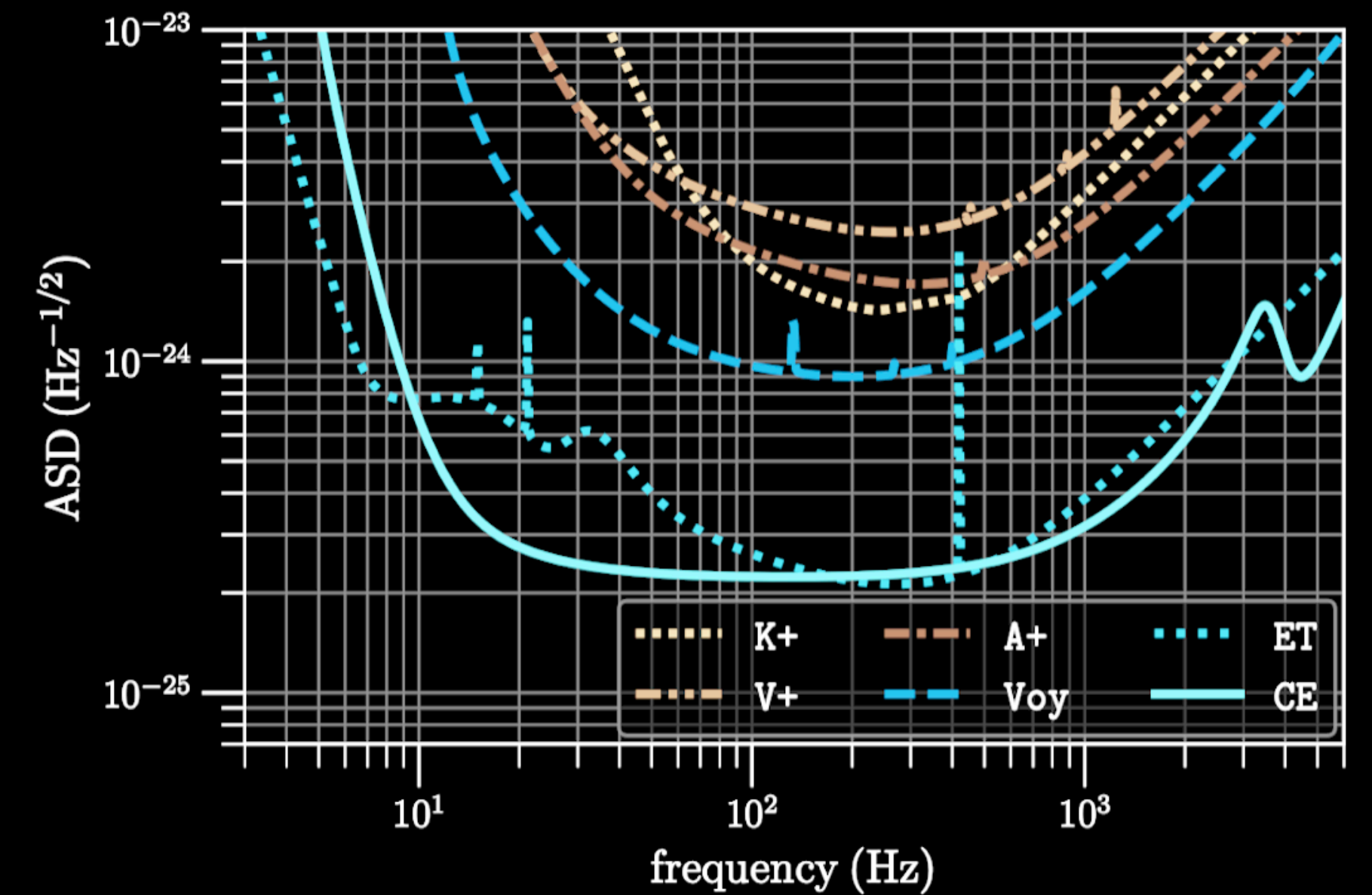
$$\epsilon(z' | \text{SNR}_{\text{thrsh}}, \sigma_{\tilde{\Lambda}, \text{thrsh}})$$

↓
Detectability of source

↓
Differentiability of source

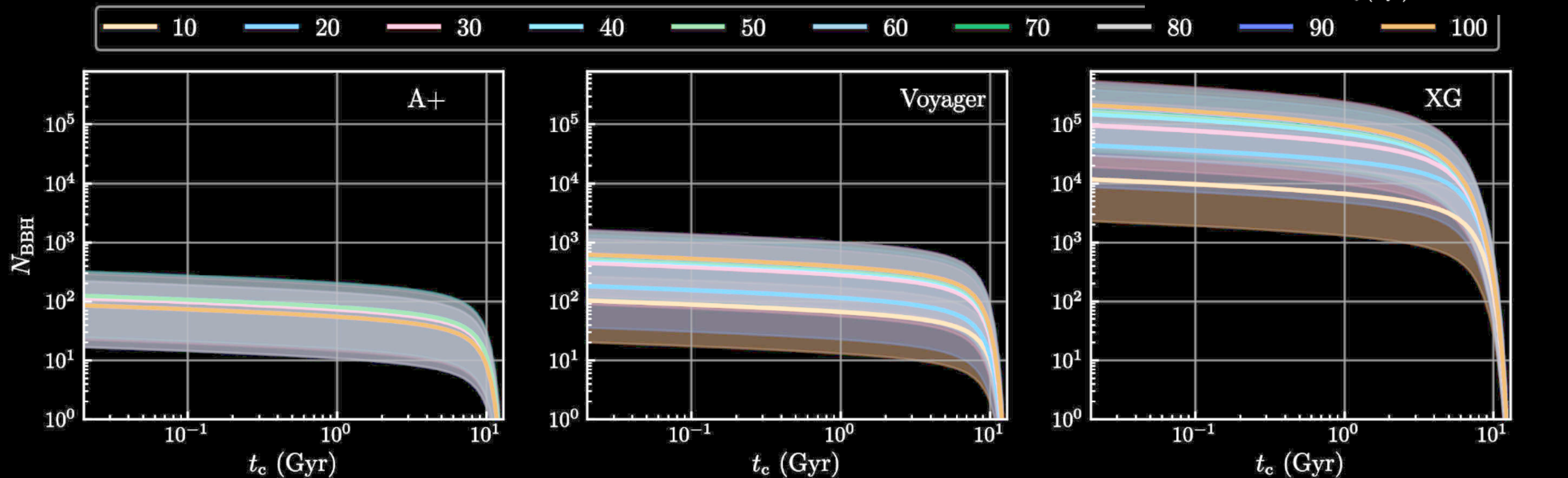
$$\tilde{\Lambda}_{\text{BBH}} = 0$$

$$\tilde{\Lambda}_{\text{BNS}} > 0$$



ASYMMETRIC DARK MATTER

BLACK HOLES FROM NEUTRON STAR IMPLOSION

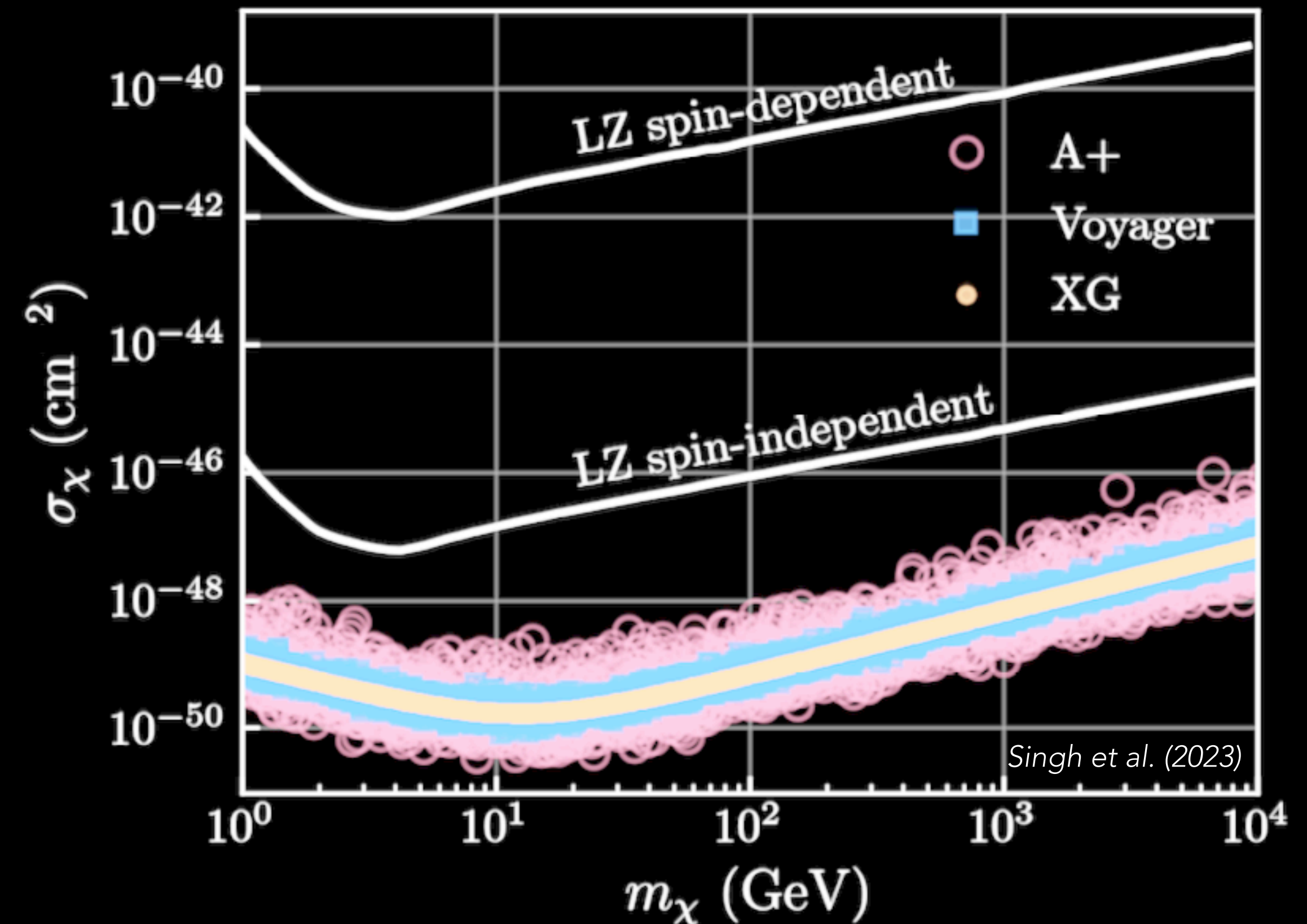


Increasingly Detectable and Differentiable

ASYMMETRIC DARK MATTER

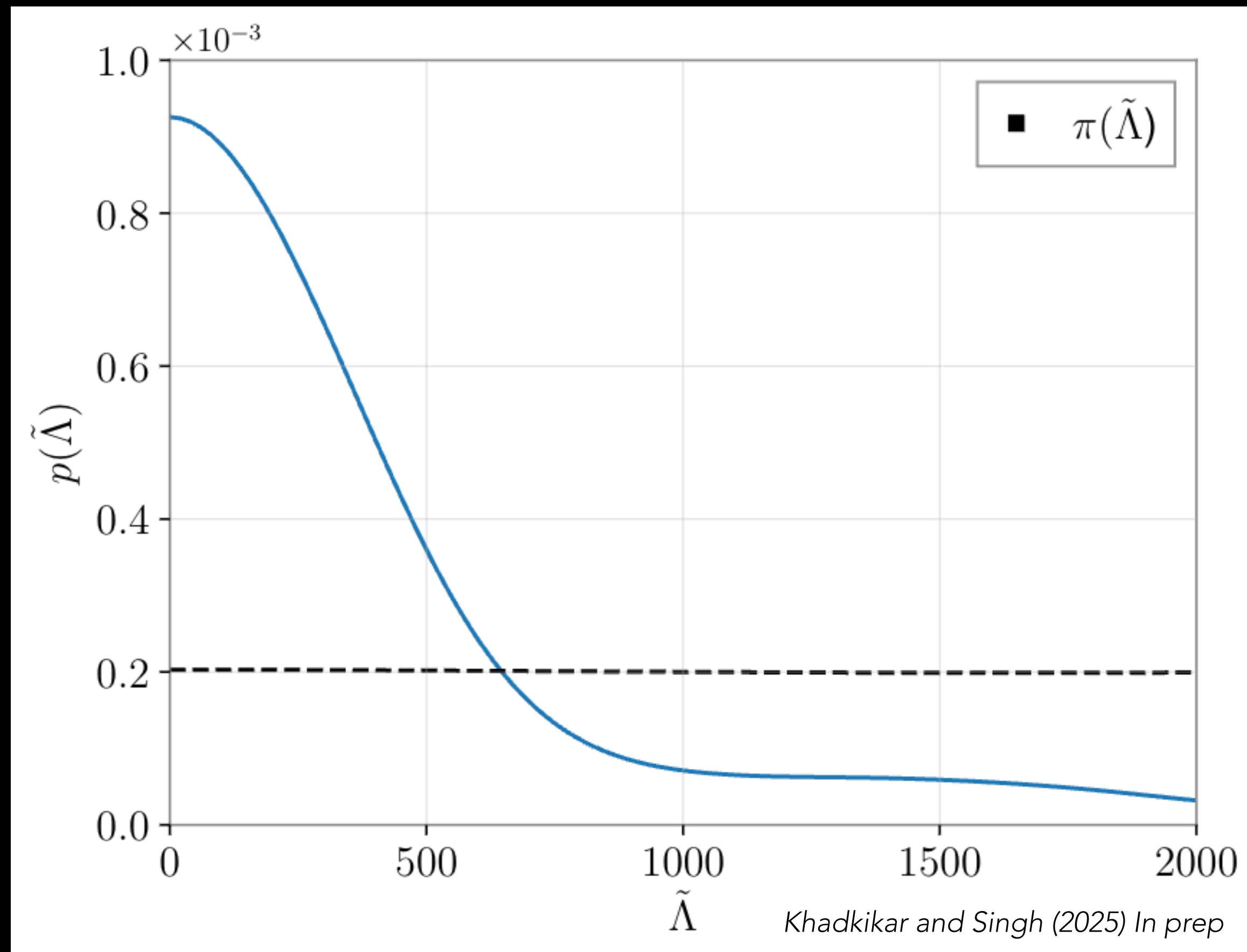
BLACK HOLES FROM NEUTRON STAR IMPLOSION

- Observed number of BBH is a function of collapse time.
- The collapse time depends on the scattering cross section and the mass of DM particle.
- Given some gravitational wave observations, we infer the collapse time and the particle properties of dark matter.



GW190425: BBH OR BNS?

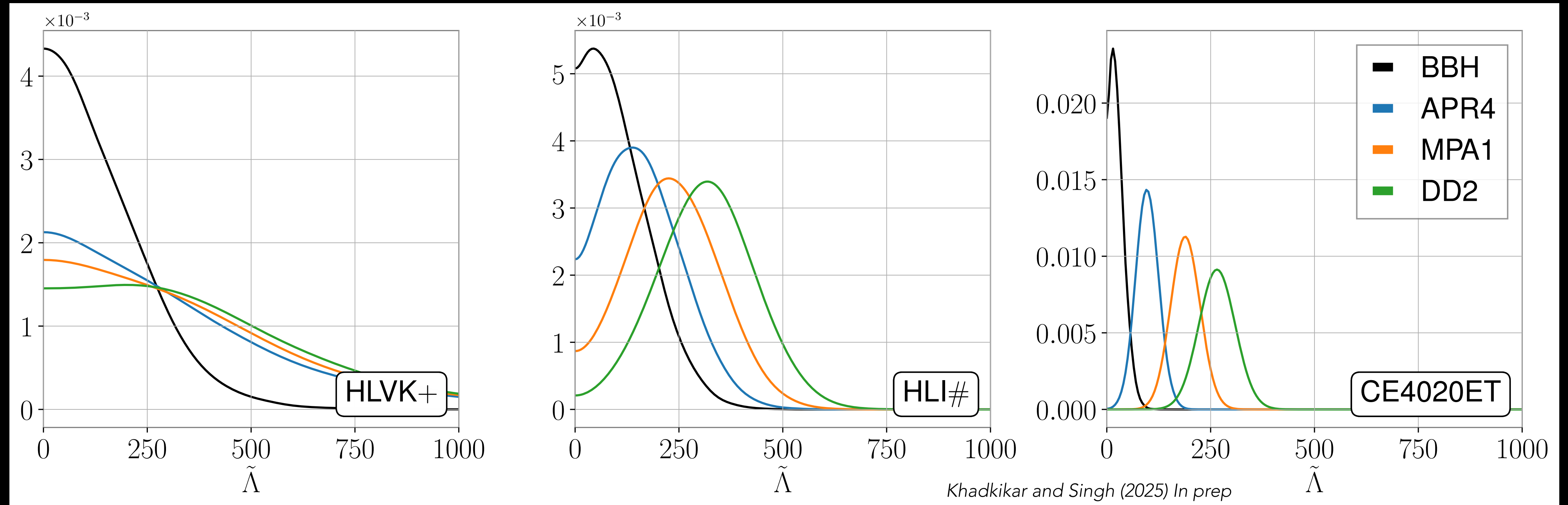
Leveraging tidal information



- The inferred effective tidal deformability, $\tilde{\Lambda}$ for GW190425
- $\tilde{\Lambda} = 0$ is not excluded.

GW190425: BBH OR BNS?

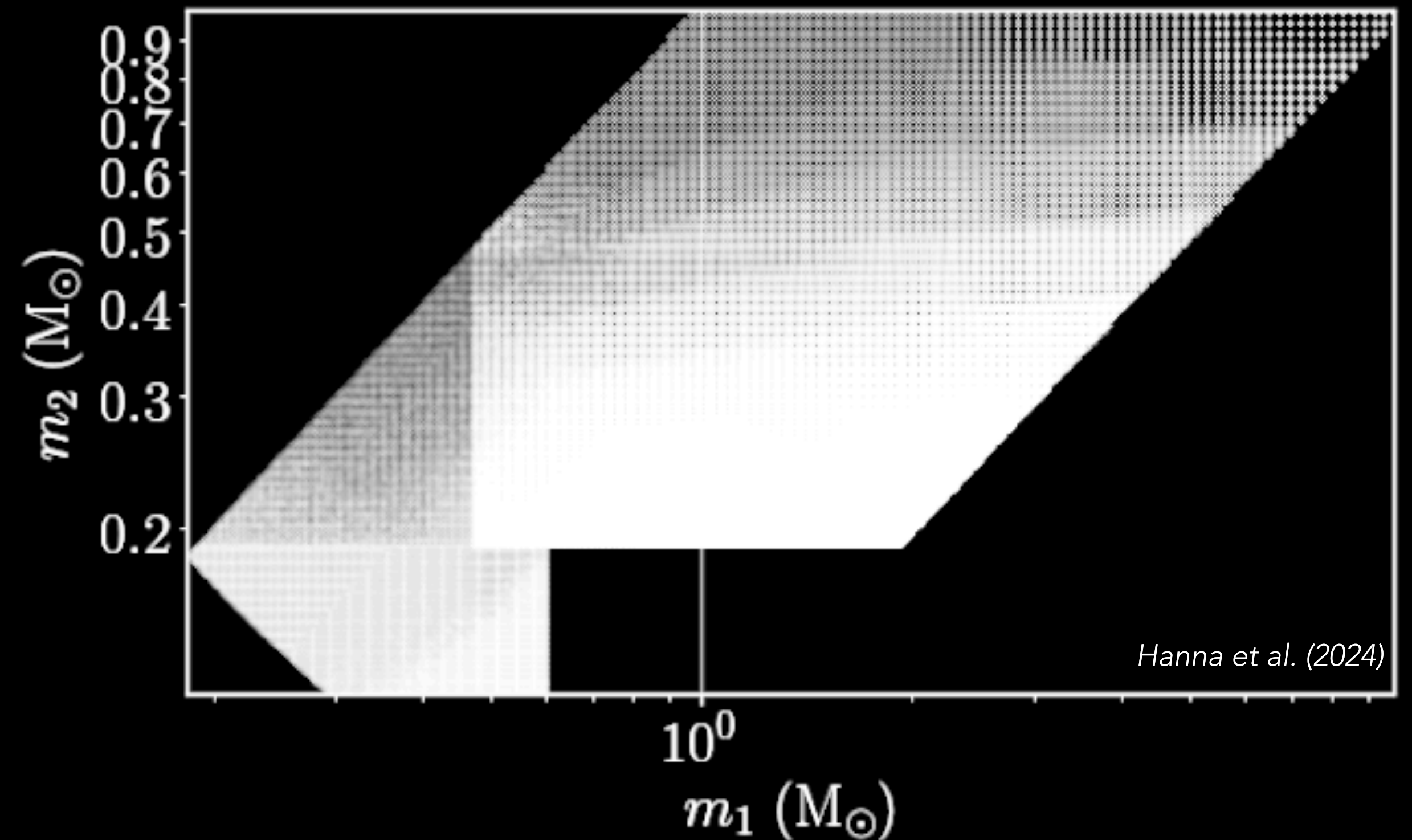
Leveraging tidal information



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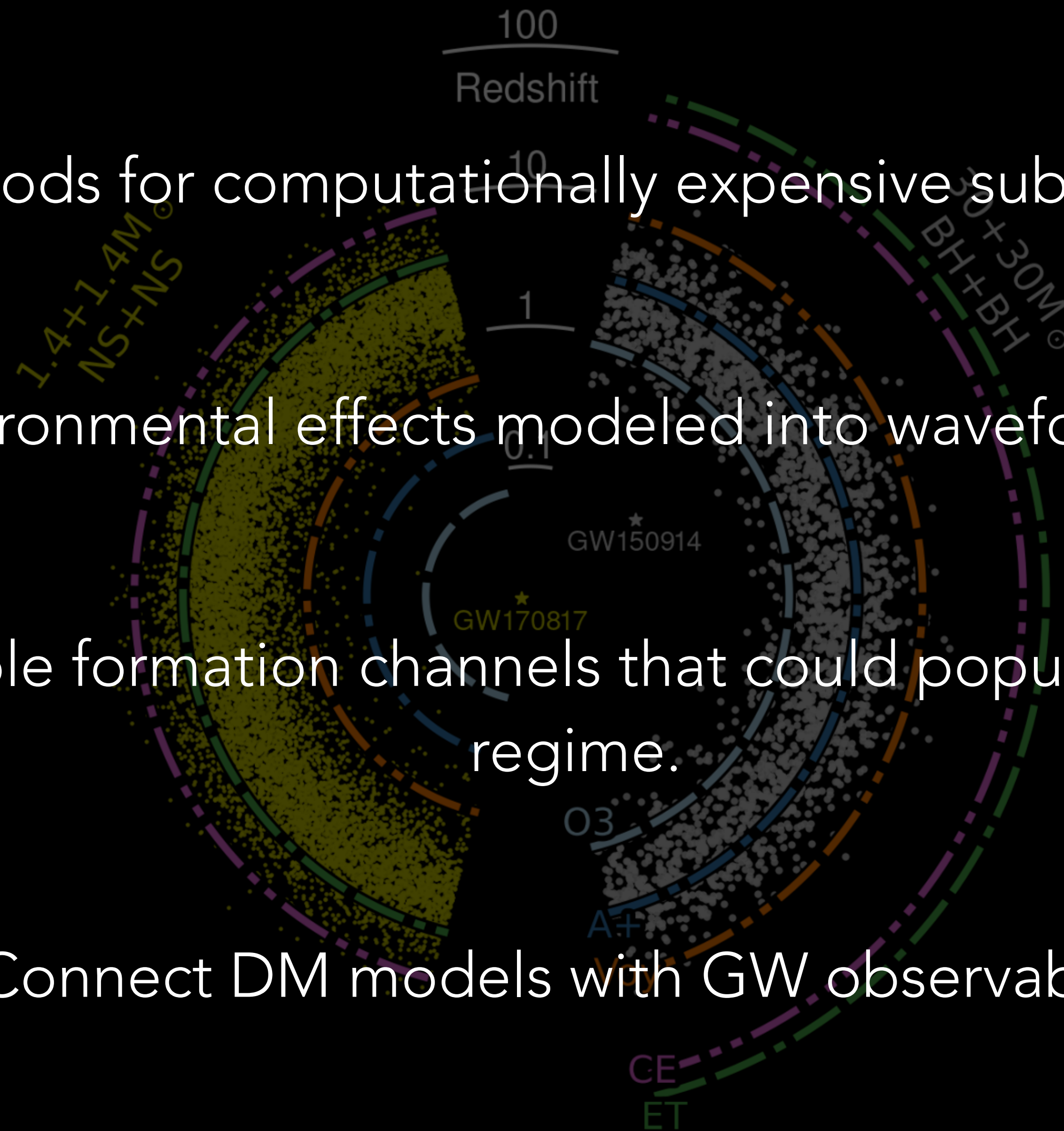
NEW METHODS FOR FUTURE DETECTORS

New search methods for computationally expensive sub-solar mass searches

Environmental effects modeled into waveforms?

Methods to decouple formation channels that could populate the sub-solar mass regime.

Connect DM models with GW observables.





Bertone et al. (2019)

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