

# Statistical Mechanics

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Here are some of the formulas I use the most in statistical mechanics

## 1 Microcanonical

The independent variables are  $E$ ,  $V$  and  $N$

### 1.1 Classical Statistics

Number of microstates

$$\Sigma \equiv \frac{1}{N!h^{3N}} \int d^{3N}q d^{3N}p \delta(E - \mathcal{H}[q_\nu, p_\nu]) \quad (1)$$

$$= \frac{1}{N!h^{3N}} \int_{E < \mathcal{H} < E + \delta E} d^{3N}q d^{3N}p \quad (2)$$

Probability Density

$$\rho[q_\nu, p_\nu] \equiv \frac{1}{\Sigma} \delta(E - \mathcal{H}[q_\nu, p_\nu]) \quad (3)$$

$$= \begin{cases} \frac{1}{\Sigma} & E < \mathcal{H}[q_\nu, p_\nu] < E + \delta E \\ 0 & \sim \end{cases} \quad (4)$$

### 1.2 Quantum Statistics

Number of microstates

$$\Sigma \equiv \sum_{\substack{l \\ E < E_l < E + \delta E}} 1 \quad (5)$$

Probability of the state  $l$

$$P_l \equiv \begin{cases} \frac{1}{\Sigma} & E < E_l < E + \delta E \\ 0 & \sim \end{cases} \quad (6)$$

Partial probability

$$P(y) \equiv \sum_{\substack{l \\ y_l=y}} P_l = \frac{1}{\Sigma} \underbrace{\sum_{\substack{l \\ y_l=y}} 1}_{\equiv w[y]} = \frac{w[y]}{\Sigma} \quad (7)$$

where  $\Sigma = \sum_l 1 = \sum_y w[y]$ . If we define the partial entropy as  $S[y] \equiv k \ln w[y]$ , then the partial probability can be written as  $P[y] \propto e^{\frac{1}{k} S[y]}$

## 2 Canonical

The independent variables are  $T$ ,  $V$  and  $N$

### 2.1 Classical Statistics

Probability density

$$\rho[q_\nu, p_\nu] \equiv \frac{e^{-\beta \mathcal{H}[q_\nu, p_\nu]}}{Z} \quad (8)$$

where  $Z \equiv \frac{1}{N! h^{3N}} \int d^{3N} q d^{3N} p e^{-\beta \mathcal{H}[q_\nu, p_\nu]}$

### 2.2 Quantum Statistics

Probability of the state  $l$

$$P_l \equiv \frac{e^{-\beta E_l}}{Z} \quad \text{where} \quad Z \equiv \sum_l e^{-\beta E_l} \quad (9)$$

Partial Probability

$$P[y] \equiv \sum_{\substack{l \\ y_l=y}} \frac{e^{-\beta E_l}}{Z} = \frac{1}{Z} \underbrace{\sum_{\substack{l \\ y_l=y}} e^{-\beta E_l}}_{\equiv Z[y]} = \frac{Z[y]}{Z} \quad (10)$$

If we define the partial Helmholtz Free Energy as  $F[y] = -k_B T \ln Z[y]$ , then the partial probability can be written as  $P[y] \propto e^{-\beta F[y]}$

### 3 Grand Canonical

The independent variables are  $T$ ,  $V$  and  $\mu$

#### 3.1 Classical Statistics

Probability density

$$\rho[N, q_\nu, p_\nu] \equiv \frac{e^{-\beta(\mathcal{H}[q_\nu, p_\nu] - \mu N)}}{\Xi} \quad (11)$$

where  $\Xi \equiv \sum_{N=0}^{\infty} \frac{1}{h^{3N}} \int d^{3N} q d^{3N} p e^{-\beta(\mathcal{H}[q_\nu, p_\nu] - \mu N)}$

#### 3.2 Quantum Statistics

Probability of the state  $l$

$$P_l \equiv \frac{e^{-\beta(E_l - \mu N_l)}}{\Xi} \quad \text{where} \quad \Xi \equiv \sum_l e^{-\beta(E_l - \mu N)} \quad (12)$$

### 4 Thermodynamics

#### 4.1 Microcanonical

Entropy:  $S = S[E, V, N] = k_B \ln \Sigma$

$$\frac{1}{T} = \frac{\partial S}{\partial E} \quad \frac{p}{T} = \frac{\partial S}{\partial V} \quad -\frac{\mu}{T} = \frac{\partial S}{\partial N} \quad (13)$$

#### 4.2 Canonical

Helmholtz Free Energy:  $F = F[T, V, N] \equiv -k_B T \ln Z = E - TS$

$$p = -\frac{\partial F}{\partial V} \quad \mu = \frac{\partial F}{\partial N} \quad S = -\frac{\partial F}{\partial T} \quad \bar{E} = -\frac{\partial \ln Z}{\partial \beta} \quad (14)$$

### 4.3 Grand Canonical

Grand canonical potential:  $\Omega = \Omega[T, V, \mu] \equiv -k_B T \ln \Xi = F - \mu N = E - TS - \mu N$

$$p = -\frac{\partial \Omega}{\partial V} \quad S = -\frac{\partial \Omega}{\partial T} \quad \bar{N} = -\frac{\partial \Omega}{\partial \mu} \quad \bar{E} - \mu \bar{N} = -\frac{\partial \ln \Xi}{\partial \beta} \quad (15)$$

## References

- [1] Diu, “Physique Statistique”
- [2] Greiner, “Thermodynamics and statistical mechanics”
- [3] Sears, “Thermodynamics, kinetic theory, and statistical thermodynamics”
- [4] Kubo, “Statistical Mechanics”